

# Odds and Ends About Osculating Walkers

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# Osculating Walkers

Consider  $p$  directed walkers on the square lattice rotated through  $45^\circ$ .

Walks take steps in the  $(1, 1)$  or  $(1, -1)$  directions.

The walkers are labelled  $k = 1, 2, \dots, p$ .

$y_t^k$  is the ordinate of the  $k$ 'th walker after  $t$  steps.

Walkers never cross but they may share vertices so  $y_t^k \leq y_t^{k+1}$ .

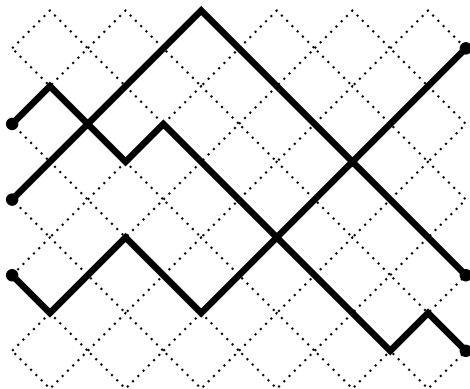
However, they are not allowed to share an edge so must separate after meeting fleetingly.

Easy to enumerate. Polynomial time algorithm of order  $n^{p-1}$  (keep track of gaps between walkers).

Generating functions solutions of Fuchsian ODEs so D-finite series.

# Star Configurations

Walkers start on neighbouring sites, but can finish anywhere.



# Known results

The dominant (physical) singularity is at  $x = x_c = 1/2^p$ .

The number of vicious stars of length  $n$  with  $p$  walkers is, from Guttmann, Owczarek and Viennot, J Phys A **31** 8123 (1998)

$$\prod_{1 \leq i \leq j \leq n} \frac{p + i + j - 1}{i + j - 1}$$

For osculating 3-stars (vicious boundaries) Mireille Bousquet-Mélou proved [J Phys Conf **42** 35 (2006)]

$$\mathcal{O}_3(x) = \frac{3 - 15x - 4x^2 - 3(1-x)\sqrt{1-8x}}{8x^2(1+x)}.$$

# Star Configurations

$p$	Init	Order	Degree	Singularities	Exponents
4	Vic	5	14	$1/16$ 4 $-1/2$	$0, 1, 2, 2, 3$ $0, 1, 2, 3, -1/2$ $0, 1, 2, -1, 1/2$
4	Osc	5	9	$1/16$ 4 $-1/2$	$0, 1, 2, 2, 3$ $0, 1, 2, 3, -1/2$ $0, 1, 2, -1, 1/2$
5	Vic	6	29	$1/32$ 2 $-1/4$	$0, 1, 2, 3, 4, 4$ $0, 1, 2, 3, -1, 1/2$ $0, 1, 2, 3, -1/2, 1/2$
5	Osc	6	12	$1/32$ 2 $-1/4$	$0, 1, 2, 3, 4, 4$ $0, 1, 2, 3, 4, 1/2$ $0, 1, 2, 3, -1/2, 1/2$

## 4 walkers – Osculating Initial Configuration

Differential operator is a direct sum of operators of order 1 and 4:

$$L_5 = L_1 \oplus L_4,$$

where  $L_4$  is a product of two order 1 operators and an order 2 operator

$$L_4 = L_2 \cdot M_1 \cdot N_1.$$

The solutions of these operators are:

$$L_1 : \frac{3 + 11x}{x(1 + 2x)}$$

$$N_1 : \frac{(1 - x)}{\sqrt{x(4 - x)} (1 + 2x)}$$

$$M_1 : \frac{(2 - 6x - 6x^2 + x^3)}{x^2 (4 - x) (1 - x) (1 + 2x)^{2/3}}$$

# 4 walkers – Osculating Initial Configuration

$$\begin{aligned} L_2 : & \frac{(10 - 135x + 24x^2 + 20x^3) \cdot \text{hypergeom}([3/2, 3/2], [3], 16x)}{2x^6 - 19x^5 + 26x^4 + 70x^3 + 10x^2 - 8x} \\ & + \frac{6x(1 - 16x)(10 - 3x - 2x^2) \cdot \text{hypergeom}([5/2, 5/2], [4], 16x)}{8x - 10x^2 - 70x^3 - 26x^4 + 19x^5 - 2x^6} \\ \\ L_2 : & \frac{(10 - 135x + 24x^2 + 20x^3) \cdot \text{hypergeom}([3/2, 3/2], [1], 1 - 16x)}{2x^6 - 19x^5 + 26x^4 + 70x^3 + 10x^2 - 8x} \\ & - \frac{18x(1 - 16x)(10 - 3x - 2x^2) \cdot \text{hypergeom}([5/2, 5/2], [2], 1 - 16x)}{8x - 10x^2 - 70x^3 - 26x^4 + 19x^5 - 2x^6} \end{aligned}$$

## 5 walkers – Osculating Initial Configuration

Differential operator is a direct sum of two operators of order 1 and and operator of order 4:

$$L_6 = L_1^a \oplus L_1^b \oplus L_4,$$

where  $L_4$  is a product of two order 1 operators and an order 2 operator

$$L_4 = L_2 \cdot M_1 \cdot N_1.$$

The solutions of these operators are:

$$L_1^a : \quad 86 - 81/x + 42/x^2$$

$$L_1^b : \quad \frac{(656386943 + 25794996300x)}{x^2\sqrt{1 + 4x}}$$

$$N_1 : \quad \frac{(2 - x)}{x^2\sqrt{1 + 4x}}$$

$$M_1 : \quad \frac{(1 + 2x - 32x^2 + 12x^3)}{x^2(1 + 4x)\sqrt{x(2 - x)}}$$



# 5 walkers – Osculating Initial Configuration

$$L_2 : \frac{(1 - 9x + 47x^2) \cdot \text{hypergeom}([-1/2, -1/2], [1], 32x)}{x^3 + 6x^4 - 24x^5 - 116x^6 + 48x^7}$$

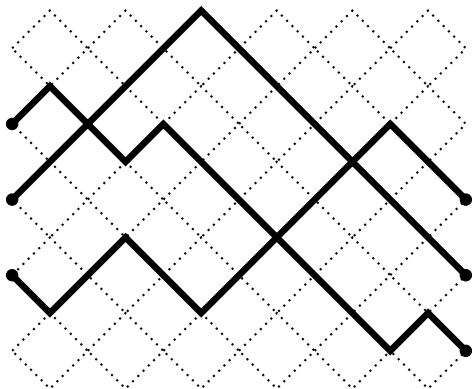
$$- \frac{2x(1 - 32x)(3 - 8x) \cdot \text{hypergeom}([1/2, 1/2], [2], 32x)}{x^3 + 6x^4 - 24x^5 - 116x^6 + 48x^7}$$

$$L_2 : \frac{(1 - 25x - 1553x^2 + 45120x^3 - 82944x^4) \cdot \text{hypergeom}([3/2, 3/2], [3], 1 - 32x)}{x^3 + 6x^4 - 24x^5 - 116x^6 + 48x^7}$$

$$- \frac{6x(1 - 32x)^3(3 - 8x) \cdot \text{hypergeom}([5/2, 5/2], [4], 1 - 32x)}{x^3 + 6x^4 - 24x^5 - 116x^6 + 48x^7}$$

# Watermelon Configurations

Walkers start and finish on neighbouring sites.



# Known results

The dominant (physical) singularity is at  $x = x_c = 1/2^p$  with critical exponent  $\alpha_p = (p^2 - 3)/2$ , and when integer valued the critical behaviour is of the form  $(1 - x/x_c)^{\alpha_p} \log(1 - x/x_c)$ .

John Essam and Tony Guttmann proved that the generating function  $\mathcal{V}_3(x)$  for vicious 3-watermelons can be expressed in terms of a Heun function

$$\begin{aligned}\mathcal{V}_3(x) &= \frac{1}{3x^3} \left[ -1 + x - 3x^2 + \text{HeunG} \left( -\frac{1}{8}, -\frac{1}{4}; -1, -2, 2, -2; -x \right) \right] \\ &= \frac{1}{3x^3} \left[ -1 + x - 3x^2 + \text{HeunG}(-8, 2; -1, -2, 2, -2; 8x) \right]\end{aligned}$$

For osculating 3-watermelons (vicious boundaries)

$$\mathcal{O}_3(x) = \frac{-1 + x + (1 - x)^2 \mathcal{V}_3(x)}{x(1 + x)}.$$

# Watermelon Configurations

$p$	Init	Order	Degree	Singularities
4	Vic	7	38	
4	Osc	7	21	$1/16, 4, -1/4, -1/2, -1$
5	Vic	10	93	
5	Osc	10	50	$1/32, 1, 2, -1/4, -1, 1 + 11x - x^2$

Simple solutions: There are 3 (4) rational solutions for  $p = 4$  (5).

For  $p = 4$  osculating boundary conditions:

$$L_7 = L_1 \oplus L_6, \quad L_6 = L_3 \cdot L_1^c \cdot L_1^b \cdot L_1^a$$

Solution of  $L_1^a$ :  $(1 - 3x^2 + 2x^3)\sqrt{4 - x}/x^{5/2}$ .

Rational solution of  $L_1^b$ .