

Mixing time of the Swendsen-Wang process on the complete graph

Tim Garoni

School of Mathematical Sciences
Monash University



Australian Government

Australian Research Council

Collaborators

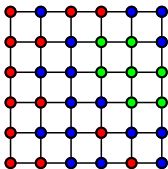
- ▶ Peter Lin (Monash University \leftrightarrow University of Washington)

Probability on Graphs

- ▶ Many problems in statistical mechanics are of the form:
 - ▶ Consider a sequence of finite graphs $G_n = (V_n, E_n)$ with:
 - ▶ $G_n \subset G_{n+1}$ and $|V_{n+1}| > |V_n|$
 - ▶ E.g. complete graphs K_n , or tori \mathbb{Z}_n^d
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 - ▶ Define (up to normalization) a probability $\pi_{n,\beta}(\cdot)$ on Ω_n

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- ▶ **Potts model:**

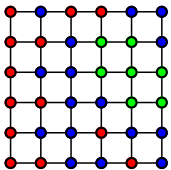


- ▶ $\Omega = [q]^V$ for fixed $q \in \{2, 3, 4, \dots\}$
- ▶ $\pi(\sigma) = \frac{1}{Z} e^{-\beta H(\sigma)}$ for $\sigma \in \Omega$
 - ▶ $H(\sigma) = -\sum_{uv \in E} \delta_{\sigma_u, \sigma_v}$
 - ▶ $\beta = 1/\text{temperature}$

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- ▶ If $\beta \approx 0$ then $\pi(\cdot) \approx$ uniform on Ω (“Disorder”)
- ▶ If $\beta \gg 1$ preference for $u \sim v$ to have $\sigma_u = \sigma_v$ (“Order”)
- ▶ Phase transition between order and disorder at critical β_c

Markov-chain Monte Carlo

- ▶ We often don't know how to normalize $\pi(\cdot)$
 - ▶ E.g. Potts partition function Z is (essentially) the Tutte polynomial

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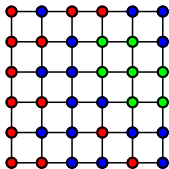
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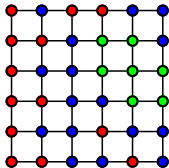
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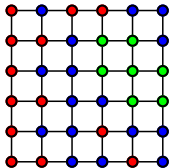
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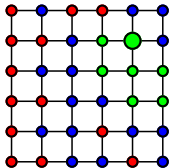
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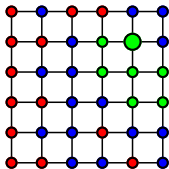
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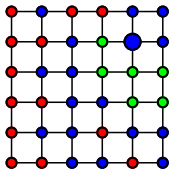


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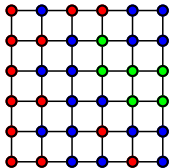


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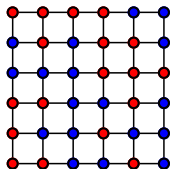
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 - ▶ If $t_{\text{mix}} = O(\text{poly}(\log |\Omega|))$ we have **rapid mixing**
 - ▶ Otherwise, we have **torpid mixing**

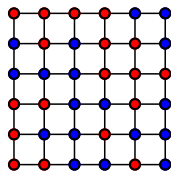
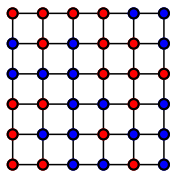
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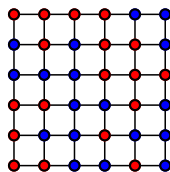
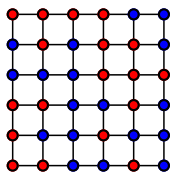
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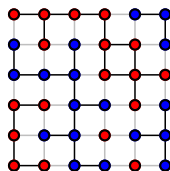
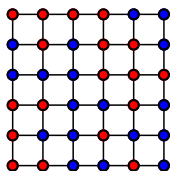


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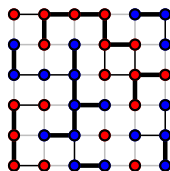
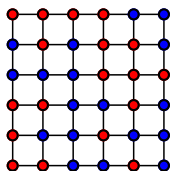


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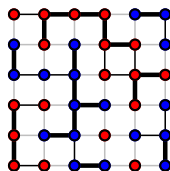
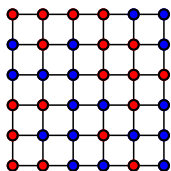


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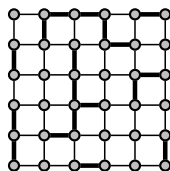
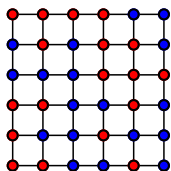


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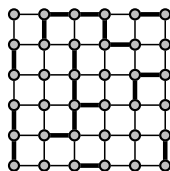
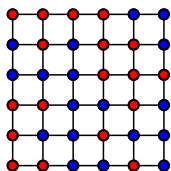


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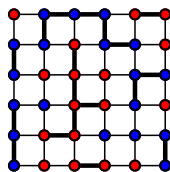
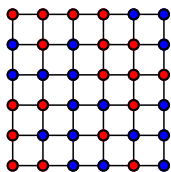


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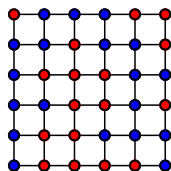
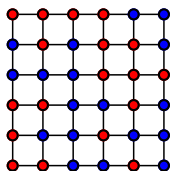


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On K_n :

- ▶ Potts model has transition at $\beta = \lambda_c/n$ with $\lambda_c = \Theta(1)$
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- ▶ Independently for each $i \in [q]$ choose Erdős-Renyi graph $\mathcal{G}(\sigma_t^{-1}(i), \lambda/n)$. Let $A_{t+1} = \cup_{i \in [q]} \mathcal{G}(\sigma_t^{-1}(i), \lambda/n)$.
- ▶ Independently and uniformly q -colour each component of (V, A_{t+1})

SW process on complete graph

On K_n :

- ▶ Potts model has transition at $\beta = \lambda_c/n$ with $\lambda_c = \Theta(1)$
 - ▶ Continuous transition for $q = 2$ (Ising)
 - ▶ Discontinuous transition for $q \geq 3$

- ▶ Potts energy depends only on magnetization $s(\sigma)$

$$-\beta H(\sigma) = \frac{\lambda}{2} n s(\sigma) \cdot s(\sigma) + \text{constant}$$

- ▶ $s^i(\sigma) = |\sigma^{-1}(i)|/n =$ fraction of vertices coloured $i \in [q]$ by σ
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Note: edge probability in $\mathcal{G}(\sigma_t^{-1}(i), \lambda/n)$ is $\lambda/n = s^i(\sigma_t)\lambda/|\sigma_t^{-1}(i)|$

Rapid mixing for $q = 2$

Potts model on K_n has **continuous** phase transition when $q = 2$

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Theorem (Cooper, Dyer, Frieze & Rue 2000)

If $q = 2$ then $SW_n(\lambda, q)$ has mixing time

$$t_{\text{mix}} = O(\sqrt{n})$$

for all $\lambda \notin (\lambda_c - \delta, \lambda_c + \delta)$ with $\delta\sqrt{\log n} \rightarrow \infty$ as $n \rightarrow \infty$.

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Theorem (Long, Nachmias, Ning, & Peres 2012)

If $q = 2$ then $\text{SW}_n(\lambda, q)$ has mixing time

$$t_{\text{mix}} = \begin{cases} \Theta(1) & \lambda < \lambda_c \\ \Theta(n^{1/4}) & \lambda = \lambda_c \\ \Theta(\log n) & \lambda > \lambda_c \end{cases}$$

- Ray, Tamayo, & Klein (1989) conjectured $n^{1/4}$ at λ_c

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Potts model on K_n has **discontinuous** phase transition when $q \geq 3$

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Potts model on K_n has **discontinuous** phase transition when $q \geq 3$

Theorem (Gore & Jerrum 1999)

If $q \geq 3$ then $\text{SW}_n(\lambda_c, q)$ has mixing time

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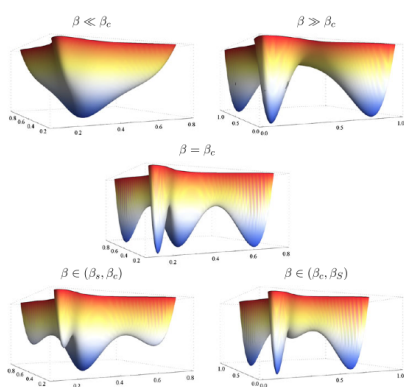
Theorem (Cuff, Ding, Loidor, Lubetzky, Peres, Sly 2012)

If $q \geq 3$ then the single-site Glauber process for the Potts model has

$$t_{\text{mix}} = \begin{cases} \Theta(n \log n) & \lambda < \lambda_o(q) \\ \Theta(n^{4/3}) & \lambda = \lambda_o(q) \\ \exp(\Omega(n)) & \lambda > \lambda_o(q) \end{cases}$$

*where $\lambda_o(q) < \lambda_c(q)$, so torpid mixing begins **before** transition*

Magnetization distribution



Large n distribution of $s(\sigma)$ known explicitly:

$$-\frac{1}{n} \log \mathbb{P}(s(\sigma) = a) \sim \phi_\lambda(a) - \inf_{a \in \Delta^{q-1}} \phi_\lambda(a)$$

$$\phi_\lambda(a) = \sum_{i=1}^q \left(a_i \log a_i - \frac{1}{2} \lambda a_i^2 \right)$$

Minima of ϕ_λ correspond either to:

- ▶ **disordered state:** $s^i = 1/q$ for all $i \in [q]$
- ▶ **ordered states:** $s^i = \alpha > 1/q$
and $s^j = \frac{1-\alpha}{q-1}$ for $j \neq i$

Figure: From Cuff *et. al* 2012

$\lambda_o(q) := \inf\{\lambda \geq 0 : \text{there exist ordered local minima of } \phi_\lambda\}$,

$\lambda_d(q) := \sup\{\lambda \geq 0 : \text{the disordered state locally minimizes } \phi_\lambda\}$.

Complete picture for $SW_n(\lambda, q)$ with $q \geq 3$

Theorem (Lin & G. 2013)

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- ▶ Gore & Jerrum's torpid mixing result extends to a non-trivial interval $(\lambda_o(q), \lambda_d(q))$ containing $\lambda_c(q)$
- ▶ Nothing special happens at $\lambda_c(q)$
- ▶ Non-trivial scaling arises at $\lambda_o(q)$
- ▶ Low and high temperature same as Ising case

Sketch of Proof

- ▶ If $Y_{t+1} := s_{t+1}^1 - \mathbb{E}[s_{t+1}^1 | \sigma_t]$ then

$$s_{t+1}^1 \approx s_t^1 + D(s_t^1) + Y_{t+1} \quad (*)$$

where

$$D_{\lambda,q}(x) := \theta(\lambda x)(1 - 1/q)x + 1/q - x$$

- ▶ $\theta(\lambda) n = \mathbb{E}(\text{size of giant component})$ in Erdős-Renyi $\mathcal{G}(n, \lambda/n)$
- ▶ $(Y_t)_{t \geq 0}$ is a sequence of martingale increments
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$$\lambda_o = \inf\{\lambda \geq 0 : D_{\lambda,q}(x) \text{ has a root on } (1/q, 1]\}$$

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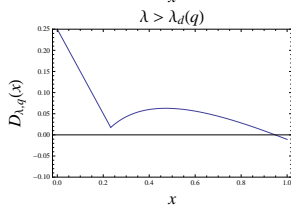
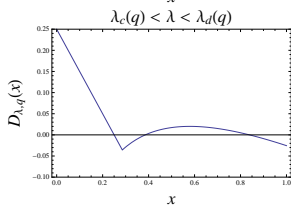
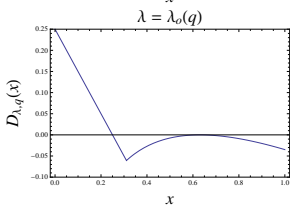
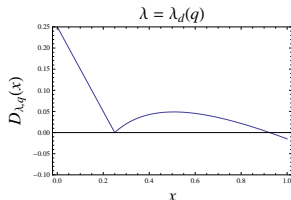
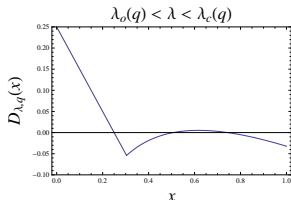
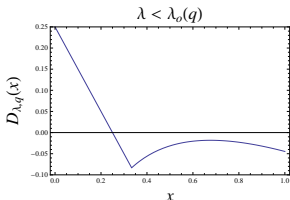
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- ▶ Coupling arguments reduce mixing time to hitting time of s_t^1

Swendsen-Wang drift

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Discussion

- ▶ Our hitting-time estimates for s_t^1 explain exponent values in mixing times for several other Potts/Ising processes
 - ▶ Mixing time exponents depend on:
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- ▶ Can one say anything for the Glauber chain for the Fortuin-Kasteleyn model?