# Lax representations of reductions of nonautonomous lattice equations

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#### Abstract

We present a method of determining a Lax representation for similarity reductions of autonomous and non-autonomous partial difference equations. We present reductions of one of the simplest integrable lattice equations, the lattice potential Korteweg-de Vries equation, which will give rise to a simple Quispel-Roberts-Thompson map, a discrete version of the first Painlevé equation and a discrete version of the fourth Painlevé equation, and their Lax pairs.

$$Q(w_{l,m}, w_{l+1,m}, w_{l,m+1}, w_{l+1,m+1}; \alpha_l, \beta_m) = 0$$
(1)



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(1)

where Q is multilinear and multidimensionally consistent and  $\alpha_I$  and  $\beta_m$  are functions of I and m respectively. Given a staircase of initial values, one evolves the system by imposing (1) on each square in  $\mathbb{Z}^2$ .



With this initial data, we may find  $w_{l,m}$  for all  $(l,m) \in \mathbb{Z}^2$ .



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$$\Psi_{l+1,m} = L_{l,m} \Psi_{l,m}$$

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$$\Psi_{l+1,m} = L_{l,m} \Psi_{l,m}$$

$$\Psi_{I,m+1} = M_{I,m}\Psi_{I,m}$$





$$\Psi_{l+1,m+1} = M_{l+1,m} L_{l,m} \Psi_{l,m}$$



$$\Psi_{l+1,m+1} = M_{l+1,m}L_{l,m}\Psi_{l,m}$$

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$$\Psi_{l+1,m+1} = M_{l+1,m}L_{l,m}\Psi_{l,m} \Longrightarrow$$
$$\Psi_{l+1,m+1} = L_{l,m+1}M_{l,m}\Psi_{l,m}$$



$$\begin{split} \Psi_{l+1,m+1} &= M_{l+1,m} L_{l,m} \Psi_{l,m} \\ \Psi_{l+1,m+1} &= L_{l,m+1} M_{l,m} \Psi_{l,m} \end{split} \Longrightarrow \qquad M_{l+1,m} L_{l,m} = L_{l,m+1} M_{l,m} \end{split}$$



is called a Lax pair if the consistency

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Implies our given lattice equation

$$Q(w_{l,m}, w_{l+1,m}, w_{l,m+1}, w_{l+1,m+1}; \alpha_l, \beta_m) = 0$$

Our toy example for this talk will be lattice KdV

$$(w_{l,m} - w_{l+1,m+1})(w_{l+1,m} - w_{l,m+1}) = \alpha_l - \beta_m$$

which has the Lax pair

$$L_{l,m} = \begin{pmatrix} w_{l,m} & \alpha_l - \gamma - w_{l,m} w_{l+1,m} \\ 1 & -w_{l+1,m} \end{pmatrix},$$
(2)  
$$M_{l,m} = \begin{pmatrix} w_{l,m} & \beta_m - \gamma - w_{l,m} w_{l,m+1} \\ 1 & -w_{l,m+1} \end{pmatrix}.$$
(3)

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Equating  $L_{I,m+1}M_{I,m}$  with  $M_{I+1,m}L_{I,m}$  gives us the above.

## Travelling Wave Reductions

Let us consider a travelling wave

## Travelling Wave Reductions



the solution satisfies a similarity constraint, in this case

$$w_{l+2,m+5}=w_{l,m},$$

so that after some time (in m), the wave has just translated in space.

## Travelling Wave Reductions



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so that after some time (in m), the wave has just translated in space. In general, we will consider reductions of the form

$$w_{l+z_1,m+z_2}=w_{l,m}.$$

which we call a  $(z_1, z_2)$ -reduction.

$$(w_{l,m} - w_{l+1,m+1})(w_{l+1,m} - w_{l,m+1}) = \alpha_l - \beta_m$$

where  $\alpha_I = \alpha = \text{const}$  and  $\beta_m = \beta = \text{const}$ , then impose a (2, 1)-reduction, then we obtain a mapping from the



We start by tracing a path going through points sharing the same value.

$$(w_{l,m} - w_{l+1,m+1})(w_{l+1,m} - w_{l,m+1}) = \alpha_l - \beta_m$$

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We label the points on the path.

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We use the equation on the lattice to define our new point

$$\hat{w}_2 = w_0 + \frac{\alpha - \beta}{w_2 - w_1}$$

$$(w_{l,m} - w_{l+1,m+1})(w_{l+1,m} - w_{l,m+1}) = \alpha_l - \beta_m$$

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So we may remove values we don't need anymore.

$$(w_{l,m} - w_{l+1,m+1})(w_{l+1,m} - w_{l,m+1}) = \alpha_l - \beta_m$$

where  $\alpha_I = \alpha = \text{const}$  and  $\beta_m = \beta = \text{const}$ , then impose a (2, 1)-reduction, then we obtain a mapping from the



We now have a new set of initial conditions that is the same as our last, but shifted in the (1,1).

$$(w_{l,m} - w_{l+1,m+1})(w_{l+1,m} - w_{l,m+1}) = \alpha_l - \beta_m$$

where  $\alpha_I = \alpha = \text{const}$  and  $\beta_m = \beta = \text{const}$ , then impose a (2, 1)-reduction, then we obtain a mapping from the



Hence, we have formed a map

$$\phi \begin{pmatrix} w_2 \\ w_1 \\ w_0 \end{pmatrix} = \begin{pmatrix} \hat{w}_2 = w_0 + \frac{\alpha - \beta}{w_1 - w_2} \\ \hat{w}_1 = w_2 \\ \hat{w}_0 = w_1 \end{pmatrix}$$

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If we allow n = 2m - l, then this is in the general form

$$w_{n+3} - w_n = \frac{\alpha - \beta}{w_{n+2} - w_{n+1}}$$

$$(w_{l,m} - w_{l+1,m+1})(w_{l+1,m} - w_{l,m+1}) = \alpha_l - \beta_m$$

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If we let  $y_n = w_{n+1} - w_n$ , this becomes

$$y_{n+1} + y_n + y_{n-1} = \frac{\alpha - \beta}{y_n}$$

which is a mapping of Quispel-Roberts-Thompson type,



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$$Y_n(\gamma) = A_n(\gamma) Y_n(\gamma)$$

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$$Y_n(\gamma) = A_n(\gamma)Y_n(\gamma)$$
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$$A_n(\gamma) \equiv M_{l+2,m}L_{l+1,m}L_{l,m}$$

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 $Y_n(\gamma) = A_n(\gamma)Y_n(\gamma)$  $Y_{n+1}(\gamma) = B_n(\gamma)Y_n(\gamma)$  $A_n(\gamma) \equiv M_{l+2,m}L_{l+1,m}L_{l,m}$  $B_n(\gamma) \equiv M_{l+1,m}L_{l,m}$ 

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Where by  $\equiv$  we mean under the variables of the reduction:



$$Y_n(\gamma) = A_n(\gamma)Y_n(\gamma)$$
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$$A_{n} = \begin{pmatrix} w_{n} & \beta - \gamma - w_{n}w_{n+2} \\ 1 & -w_{n+2} \end{pmatrix} \begin{pmatrix} w_{n+1} & \alpha - \gamma - w_{n}w_{n+1} \\ 1 & -w_{n} \end{pmatrix}$$
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The consistency relation for these two equations is equivalent to



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$$Y_{n+1} = A_{n+1}B_nY_n$$
$$Y_{n+1} = B_nA_nY_n$$
$$A_{n+1}B_n = B_nA_n$$
$$w_{n+3} - w_n = \frac{\alpha - \beta}{w_{n+2} - w_{n+1}}$$

Note: We could express this in terms of the  $y_n$  (with some effort).

$$A_{n} = \begin{pmatrix} w_{n} & \beta - \gamma - w_{n}w_{n+2} \\ 1 & -w_{n+2} \end{pmatrix} \begin{pmatrix} w_{n+1} & \alpha - \gamma - w_{n}w_{n+1} \\ 1 & -w_{n} \end{pmatrix}$$
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## Deautonomizing the theory

If we impose

$$w_{l+z_1,m+z_2}=w_{l,m},$$

we find a constraint that the equation describing the evolution at (I, m) must be the same as the equation at  $(I + z_1, m + z_2)$ . In our toy example, imposing the above gives

$$(w_{l,m} - w_{l+1,m+1})(w_{l+1,m} - w_{l,m+1}) - \alpha_l + \beta_m = 0,$$
  
$$(w_{l,m} - w_{l+1,m+1})(w_{l+1,m} - w_{l,m+1}) - \alpha_{l+z_1} + \beta_{m+z_2} = 0.$$

is consistent when

$$\alpha_{I+z_1} - \alpha_I = \beta_{m+z_2} - \beta_m := \operatorname{lcm}(z_1, z_2)h$$

where h is constant in l and m by a separation of variables argument.

In our example, the (2, 1)-reduction, we have

$$\alpha_{I+2} - \alpha_I = \beta_{m+1} - \beta_m = 2h$$

The simplest solution is where these are just linear;

$$\alpha_I = lh + a, \qquad \beta_m = mh + b.$$

Our reduction is then (remembering n = 2m - l)

$$w_{n+3} - w_n = \frac{\alpha_{l+1} - \beta_m}{w_{n+2} - w_{n+1}} = \frac{a - b - hn + h}{w_{n+2} - w_{n+1}}$$

which by letting  $y_n = w_{n+1} - w_n$  we obtain the equation

$$y_{n+1} + y_n + y_{n-1} = \frac{nh + b - a}{y_n},$$

which is a discrete version of Painlevé I.

To form the Lax representation here, we find in the original Lax pair, the parameters,  $\alpha_l$  and  $\beta_m$ , always appear with  $\gamma$ ,

$$L_{l,m} = \begin{pmatrix} w_{l,m} & \alpha_l - \gamma - w_{l,m}w_{l+1,m} \\ 1 & -w_{l+1,m} \end{pmatrix},$$
(4)  
$$M_{l,m} = \begin{pmatrix} w_{l,m} & \beta_m - \gamma - w_{l,m}w_{l,m+1} \\ 1 & -w_{l,m+1} \end{pmatrix}.$$
(5)

The key part is that we introduce a spectral parameter that couples a direction in (I, m)-space with  $\gamma$ , we let

$$x = lh - \gamma.$$

now the Lax pairs can be written in terms of x and n, as

$$\alpha_{I} - \gamma = \mathbf{a} + \underbrace{\mathbf{hI} - \gamma}_{\mathbf{x}}$$
$$\beta_{m} - \gamma = \mathbf{b} + \underbrace{\beta_{m} - \alpha_{I}}_{n\mathbf{h}} + \underbrace{\mathbf{hI} - \gamma}_{\mathbf{x}}$$



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$$Y_n(x+2h) = A_n(x)Y_n(x)$$
$$Y_{n+1}(x+h) = B_n(x)Y_n(x)$$
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Where we need to now take into account the non-autonomy:

$$A_{n} = \begin{pmatrix} w_{n} & b + nh + x - w_{n}w_{n+2} \\ 1 & -w_{n+2} \end{pmatrix} \begin{pmatrix} w_{n+1} & x + a + h - w_{n}w_{n+1} \\ 1 & -w_{n} \end{pmatrix}$$
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If we examine our original condition,

$$\alpha_{l+2} - \alpha_l = \beta_{m+1} - \beta_m = 2h,$$

a little more closely, we find that we may build an extra variable. The general solution is

$$\alpha_{I} = \begin{cases} hI + a_{1} & \text{where } I \text{ is odd,} \\ hI + a_{2} & \text{where } I \text{ is even,} \end{cases}$$

and  $\beta_m = 2mh + b$ , but since the resulting equation only depends on  $\alpha_l - \beta_m$ , b can be chosen to be 0.

The single shift, 
$$n \rightarrow n + 1$$
 has the effect  
 $w_{n+2} = w_{n+1}$   $(w_n - w_{n+3})(w_{n+1} - w_{n+2}) = hn + a_2.$   
 $a_1 \rightarrow a_2 + h, \quad a_2 \rightarrow a_1 - h.$   
But the double shift has the desirable effect

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$$(w_n - w_{n+3})(w_{n+1} - w_{n+2}) = hn + a_2,$$
  
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 $(w_{n+1} - w_{n+4})(w_{n+2} - w_{n+3}) = hn + a_1,$ 

However, if we let

$$y_n = (w_{n+2} - w_n)(w_{n+2} - w_{n+1}) - a_2,$$
  
$$z_n = \frac{(w_{n+2} - w_n)(y_n + a_1)}{y_n + nh},$$

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$$(w_n - w_{n+3})(w_{n+1} - w_{n+2}) = hn + a_2,$$
  
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then the system is equivalent to

$$y_n + y_{n+2} = z_n^2 - (a_1 + a_2),$$
  
$$z_n z_{n+2} = -\frac{(y_{n+2} + a_1)(y_{n+2} + a_2)}{(y_{n+2} + (n+2)h)}.$$

$$(w_n - w_{n+3})(w_{n+1} - w_{n+2}) = hn + a_2,$$
  
 $(w_{n+1} - w_{n+4})(w_{n+2} - w_{n+3}) = hn + a_1,$ 

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then the system is equivalent to

$$(y + \tilde{y} + a_1 + a_2)(y + \underline{y} + a_1 + a_2) = \frac{(y + a_1)^2(y + a_2)^2}{(y + nh)^2}$$

where  $y_n = y$ ,  $y_{n+2} = \tilde{y}$ ,  $y_{n-2} = \tilde{y}$  and  $\tilde{n} = n + 2$ . This is a discrete version of the fourth Painlevé equation.

The Lax pair in the  $y_n$  and  $z_n$  variables are given by

$$Y_n(x+2h) = A_n(x)Y_n(x),$$
  
$$Y_{n+2}(x) = B_n(x)Y_n(x),$$

where

$$A_n(x) \equiv \begin{pmatrix} -\frac{(y_n+a_1)(y_n+a_2)}{z_n} & x^2+\delta x+\epsilon\\ x-y_n & (y_n+nh)z_n \end{pmatrix},$$

where

$$\delta = y_n + a_1 + a_2 + hn, \epsilon = (y_n + hn)(y_n + a_1) + (hn + y_n + a_1)a_2,$$

and

$$B_n(x) = \begin{pmatrix} -z_n & x + hn + z_n^2 \\ 1 & -z_n \end{pmatrix}.$$

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We have also performed a reduction from the lattice modified Korteweg-de Vries equation

$$\alpha_{l}(w_{l,m}w_{l+1,m}-w_{l,m+1}w_{l+1,m+1})-\beta_{m}(w_{l,m}w_{l,m+1}-w_{l+1,m}w_{l+1,m+1})=0$$

to a q-analogue of the sixth Painlevé equation

$$\begin{split} y \hat{y} &= \frac{q^2 \left(q^2 b_1 t^2 + z a_2\right) \left(b_2 t^2 + z a_1\right)}{\left(z b_1 q^2 + a_2\right) \left(a_1 + z b_2\right)},\\ z \hat{z} &= \frac{\left(t^2 b_1 q^4 + \hat{y} a_1\right) \left(t^2 b_2 q^4 + \hat{y} a_2\right)}{q^2 \left(a_1 + \hat{y} b_1\right) \left(a_2 + \hat{y} b_2\right)}, \end{split}$$

where  $\hat{t} = q^2 t$  for some fixed  $q \in \mathbb{C}$  and the  $a_i$  and  $b_j$  are fixed parameters.

A reduction from the lattice Schwarzian Korteweg-de Vries equation

$$\alpha_{l} \left( \frac{1}{w_{l,m+1} - w_{l+1,m+1}} + \frac{1}{w_{l+1,m} - w_{l,m}} \right)$$

$$= \beta_{m} \left( \frac{1}{w_{l+1,m+1} - w_{l+1,m}} + \frac{1}{w_{l,m} - w_{l,m+1}} \right),$$
(6)

to a q-analogue of the sixth Painlevé equation

$$y\hat{y} = \frac{b_1b_2(z-1)(q^2z-1)}{q^2(b_1b_2t - \theta_1 z)(b_1b_2t - \theta_2 z)},$$

$$z\hat{z} = \frac{(b_1q^2t\hat{y} - 1)(b_2q^2t\hat{y} - 1)}{q^2(a_1\hat{y} - 1)(a_2\hat{y} - 1)},$$
(8)

where  $\hat{t} = q^2 t$  for some fixed  $q \in \mathbb{C}$  and the  $a_i$ ,  $b_j$  and  $\theta_i$  are fixed parameters.

Another reduction from the lattice Schwarzian Korteweg-de Vries equation

$$\alpha_{l} \left( \frac{1}{w_{l,m+1} - w_{l+1,m+1}} + \frac{1}{w_{l+1,m} - w_{l,m}} \right)$$

$$= \beta_{m} \left( \frac{1}{w_{l+1,m+1} - w_{l+1,m}} + \frac{1}{w_{l,m} - w_{l,m+1}} \right),$$
(9)

to an equation just above the previous equation, known as q-P( $E_6^{(1)}$ ):

$$\begin{aligned} (\hat{y}z-1)(\hat{y}\hat{z}-1) &= \frac{(a_1\hat{y}-1)(a_2\hat{y}-1)(a_3\hat{y}-1)(a_4\hat{y}-1)}{(b_1q^4t\hat{y}-1)(b_2q^4t\hat{y}-1)},\\ (yz-1)(\hat{y}z-1) &= \frac{\theta_1(z-a_1)(z-a_2)(z-a_3)(z-a_4)}{(b_1b_2tz+\theta_1)(a_1a_2a_3a_4+\theta_1q^4tz)}, \end{aligned}$$

where  $\hat{t} = q^4 t$ , the  $a_i$ ,  $b_i$  and  $\theta_1$  are fixed parameters and  $q \in \mathbb{C}$ .

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