

# Parallel Tempering and Population Annealing for Rough Free Energy Landscapes

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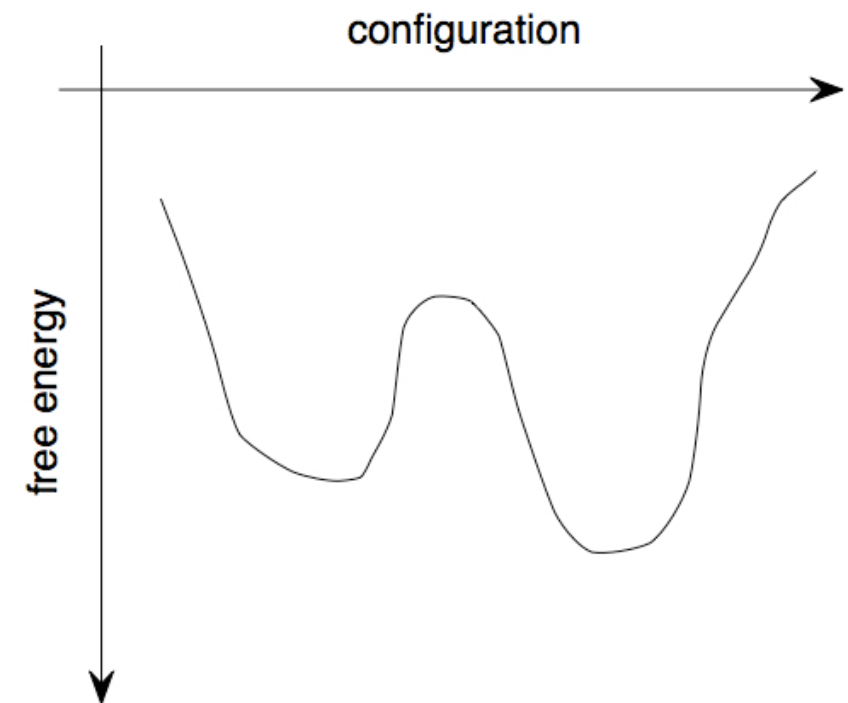
Melbourne, July 26, 2010

# Collaborators

- Stephan Burkhardt
- Burcu Yucesoy
- Helmut Katzgraber

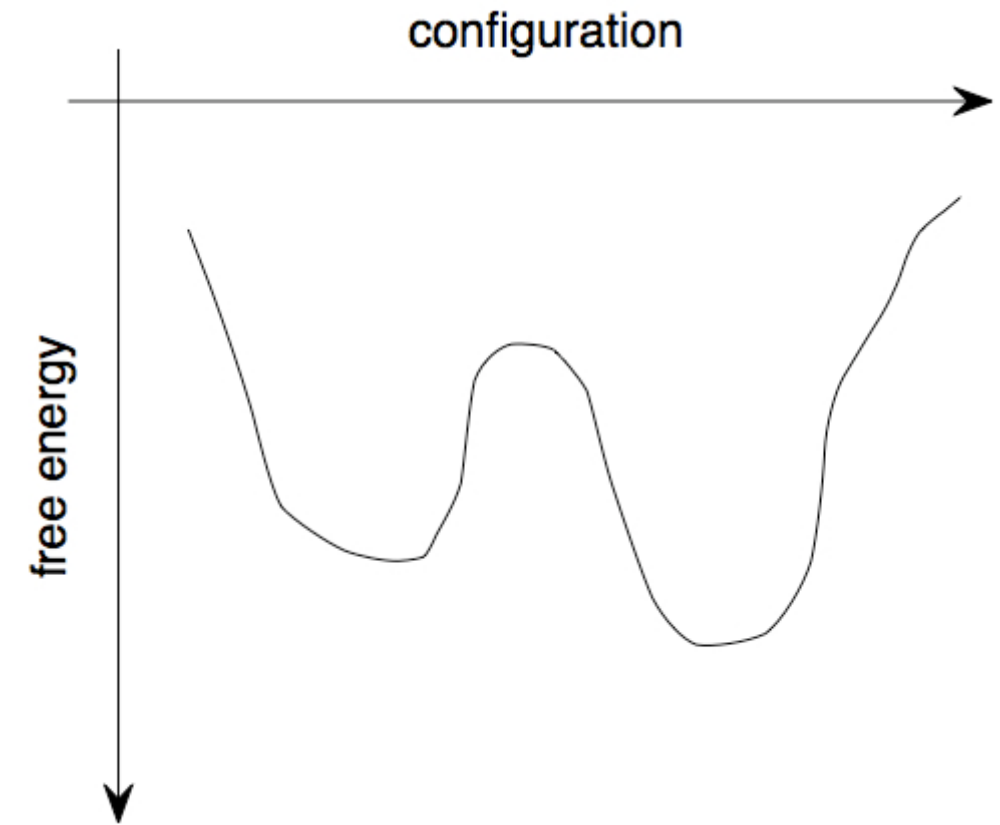
# Motivation

- How to sample equilibrium states of systems with rugged free energy landscapes, e.g. spin glasses, configurational glasses, proteins, NP-hard combinatorial optimization problems.



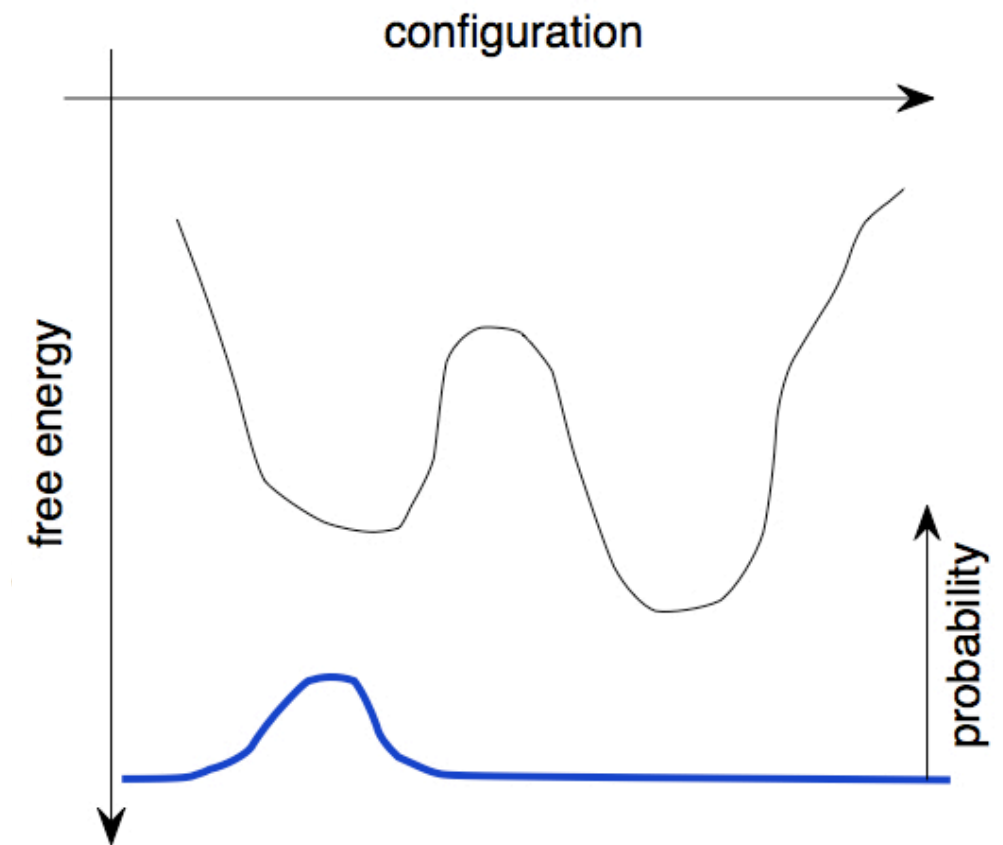
# Problem

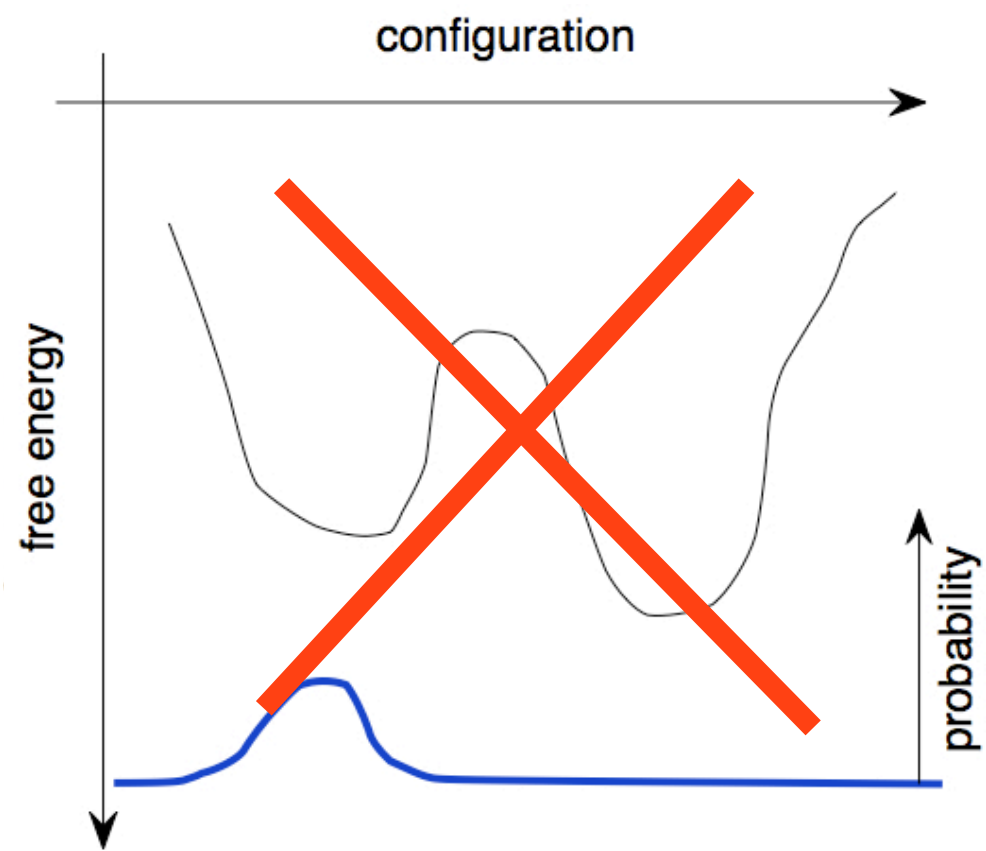
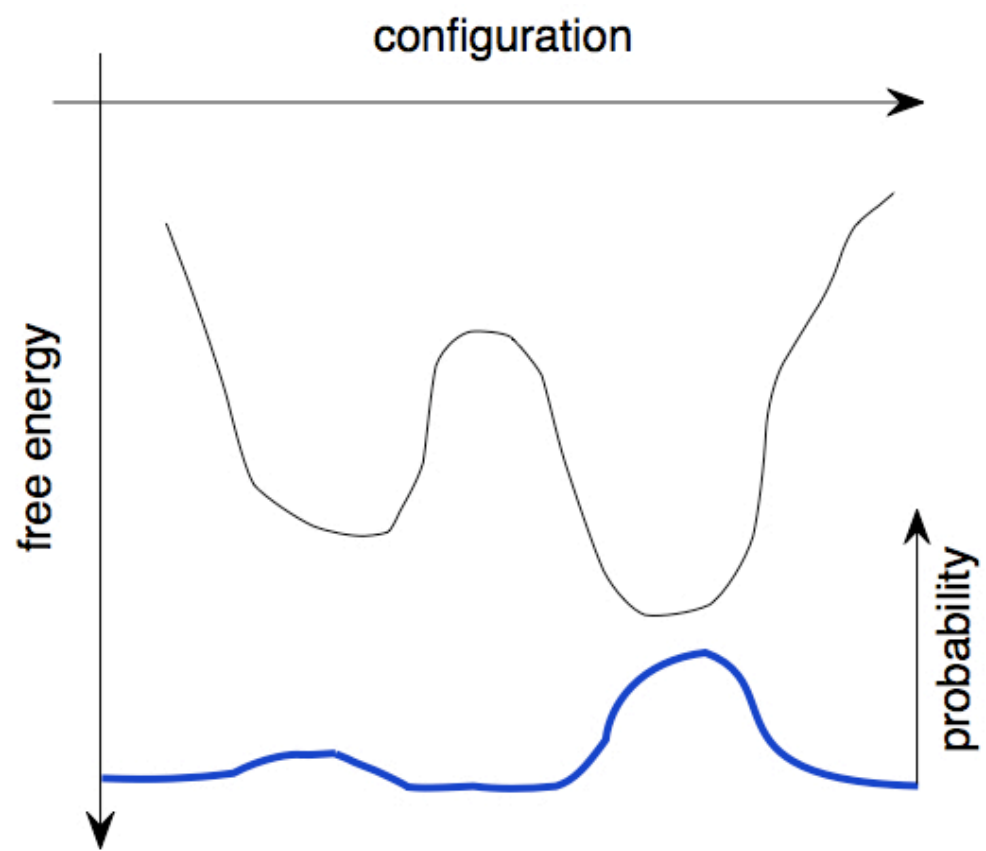
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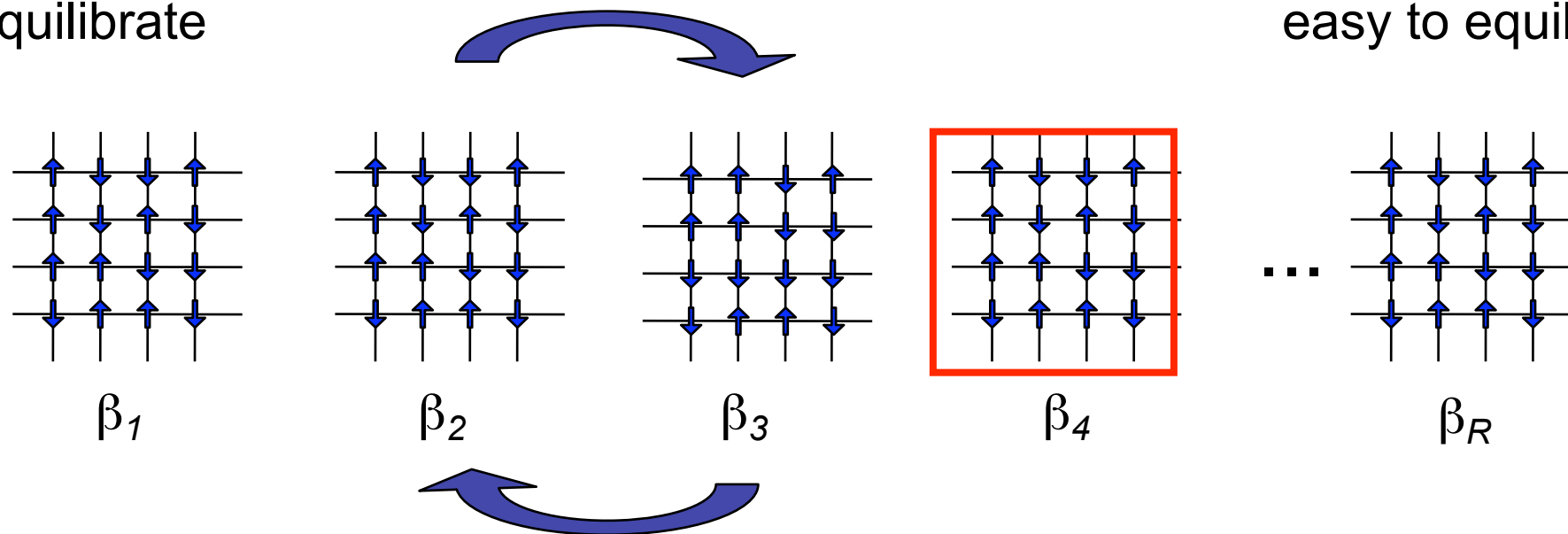
# Outline

- Introduction to Parallel Tempering (aka Replica Exchange Monte Carlo)
- Strengths and Weaknesses of PT (in several simple landscapes)
- Introduction to Population Annealing
- PA vs PT
- Conclusions

# Parallel Tempering

hard to equilibrate

easy to equilibrate



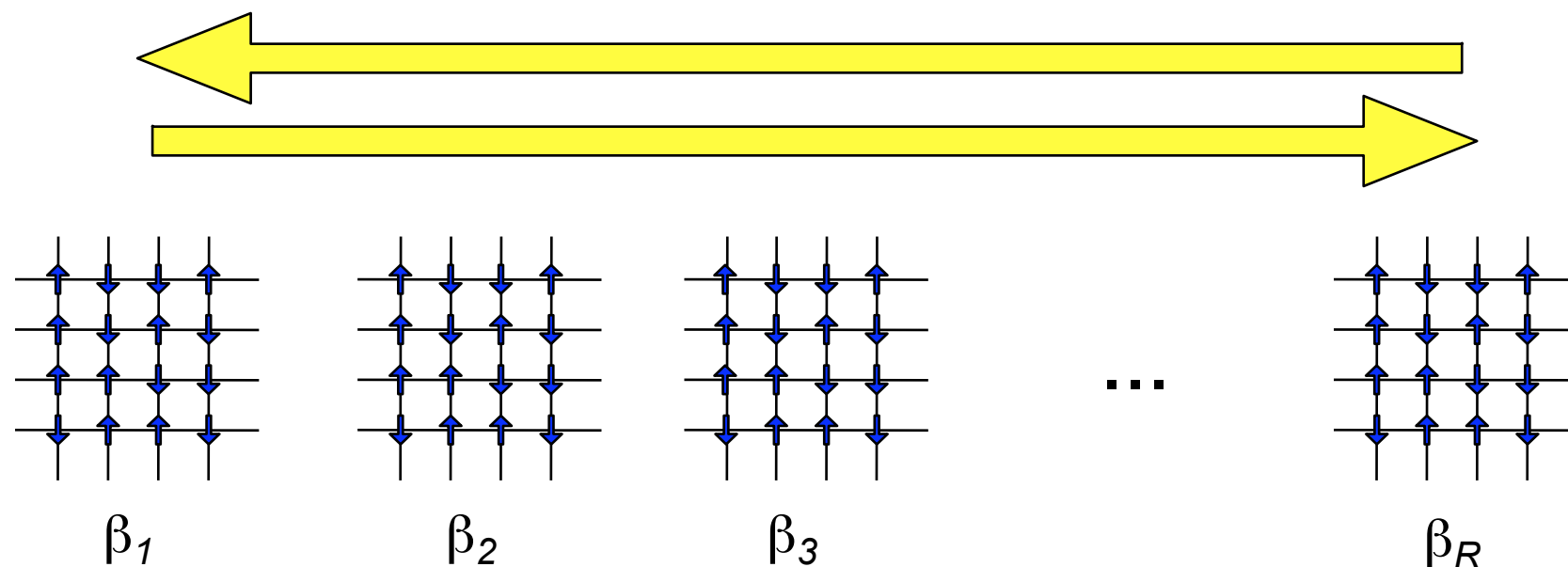
- $R$  replicas at inverse temperatures  $\beta_1 > \beta_2 > \dots > \beta_R$  (each with the same couplings).
- MCMC (e.g. Metropolis) on each replica
- Exchange replicas with energies  $E$  and  $E'$  and temperatures  $\beta$  and  $\beta'$ , with probability:

$$p_{\text{swap}} = \min \left[ 1, e^{(\beta - \beta')(E - E')} \right]$$



# Intuition

- Mixing is accelerated by “round trips” from low to high temperature and back.

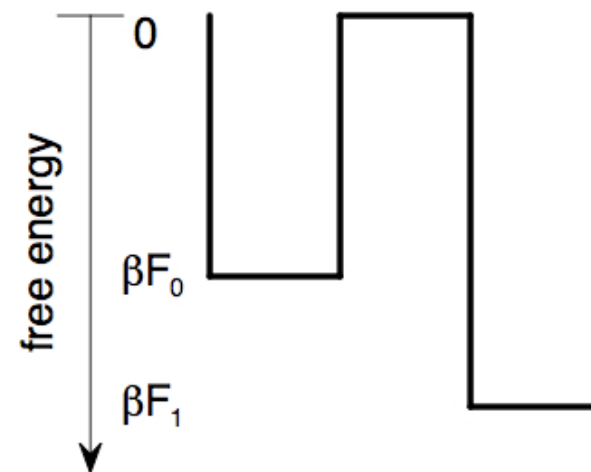


- Parallel tempering can be optimized by minimizing the equilibrium round trip time.

# A simple landscape:

- Consider a model free energy landscape with two free energy minima separated by a high barrier.

– JM, PRE **80**, 056706 (2009)



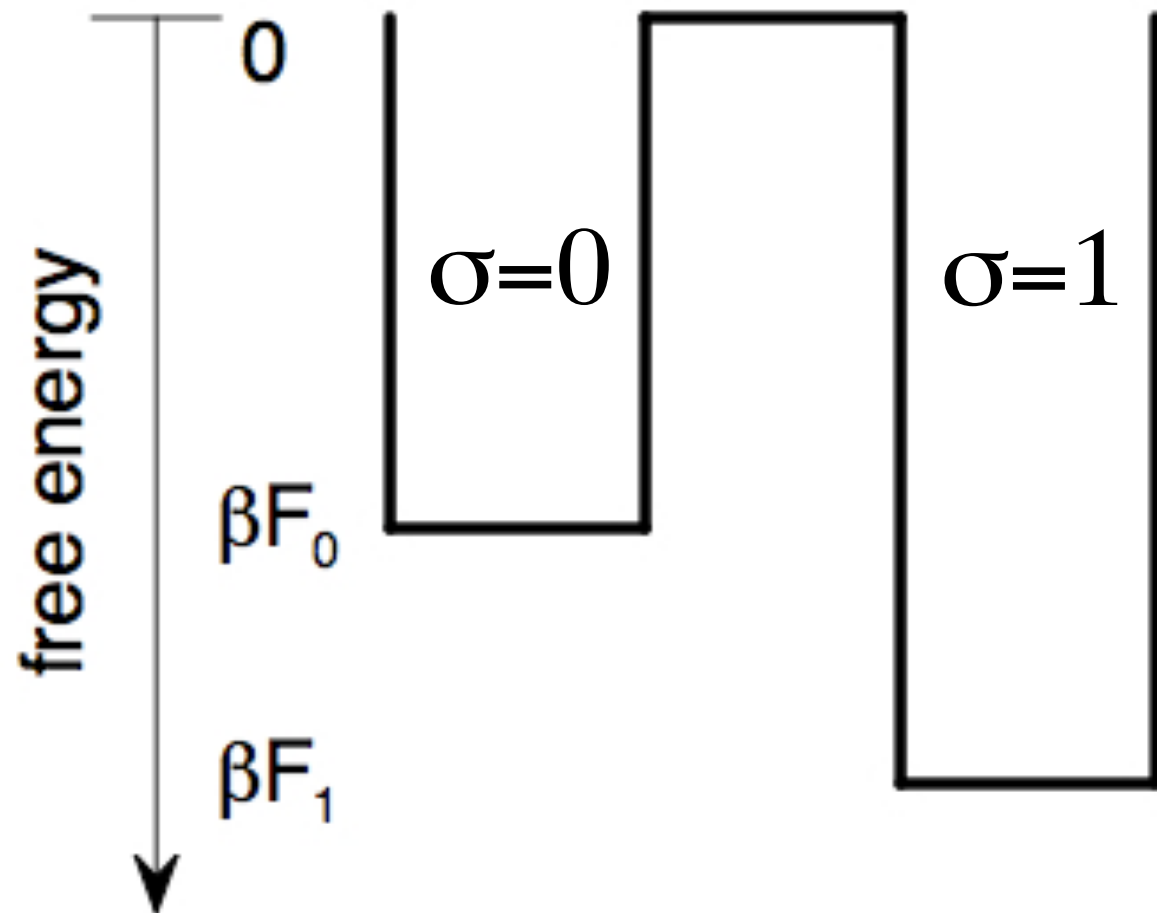
# Two Well Model

$$\beta F_{\sigma}(\beta) = -\frac{1}{2}(\beta - \beta_c)^2(K + H\sigma)$$

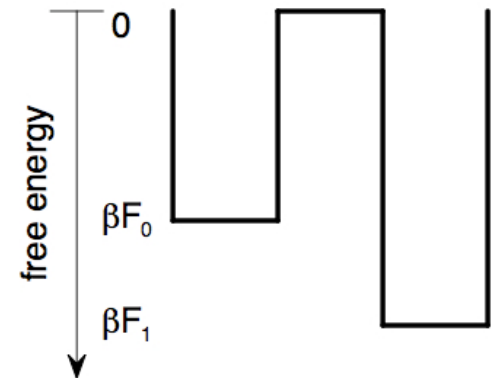
$$\sigma = 0, 1$$

$$\beta \delta F = -\frac{1}{2}(\beta - \beta_c)^2 H$$

$$\text{Prob}[\sigma = +1] = \frac{1}{1 + e^{-\beta \delta F}}$$



# Parallel tempering for the two-well model

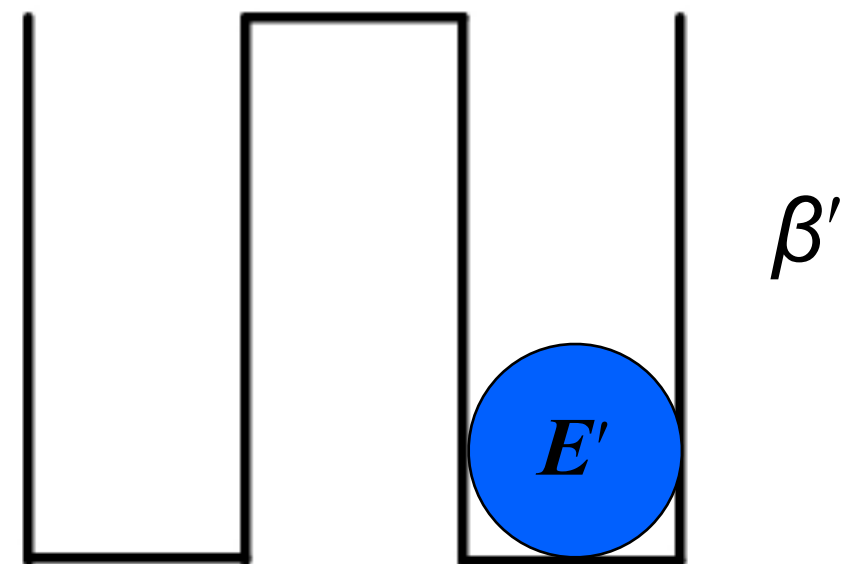
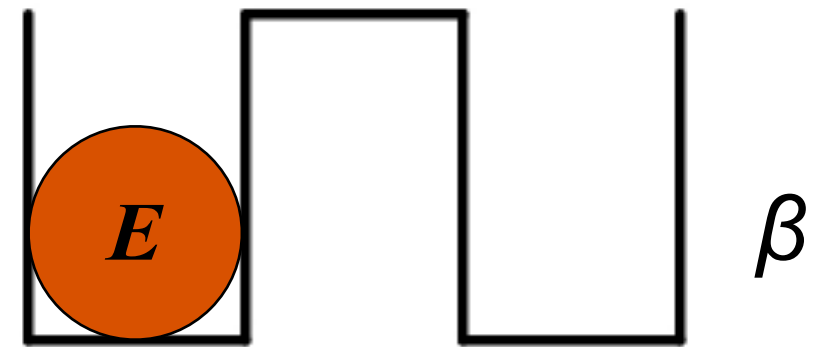


- Assumptions:
  - Fast equilibration within each well by standard MCMC.
  - No transitions between wells except at  $\beta_c$  where each well is equally probable.
  - Energy is normally distributed in each state; from thermodynamics:

$$\langle E \rangle = -(\beta - \beta_c)(K + H\sigma) \quad \mathbf{Var}(E) = (K + H\sigma)$$

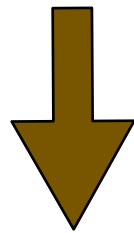
# Replica exchange probabilities

- For the two-well model, replica exchange transition probabilities can be computed exactly. For symmetric wells ( $H=0$ ):



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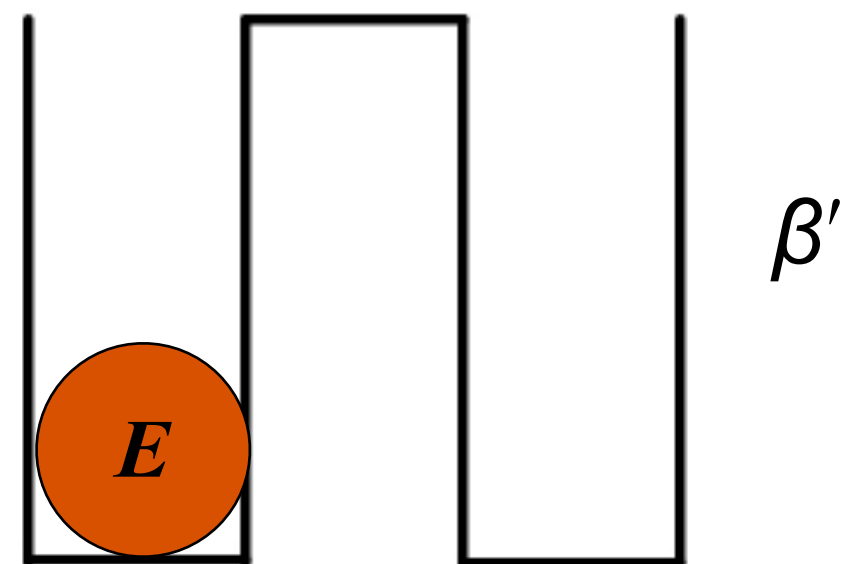
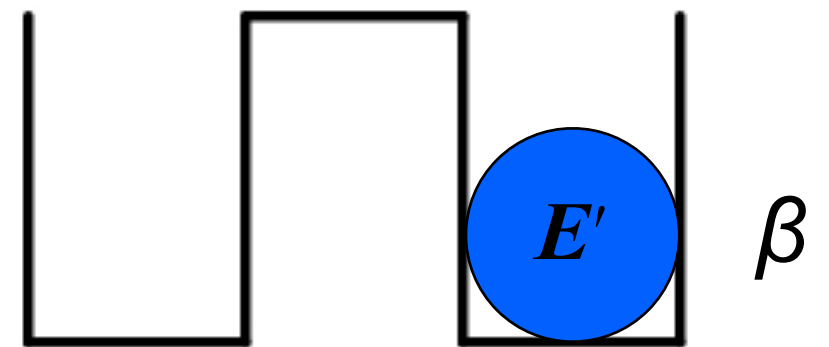
$$\langle E \rangle = -(\beta - \beta_c)K, \quad \text{Var}(E) = K$$



$$p_{\text{swap}} = \frac{1}{2} \text{Erfc} \left( (\beta - \beta') \sqrt{K} \right)$$

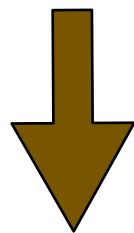
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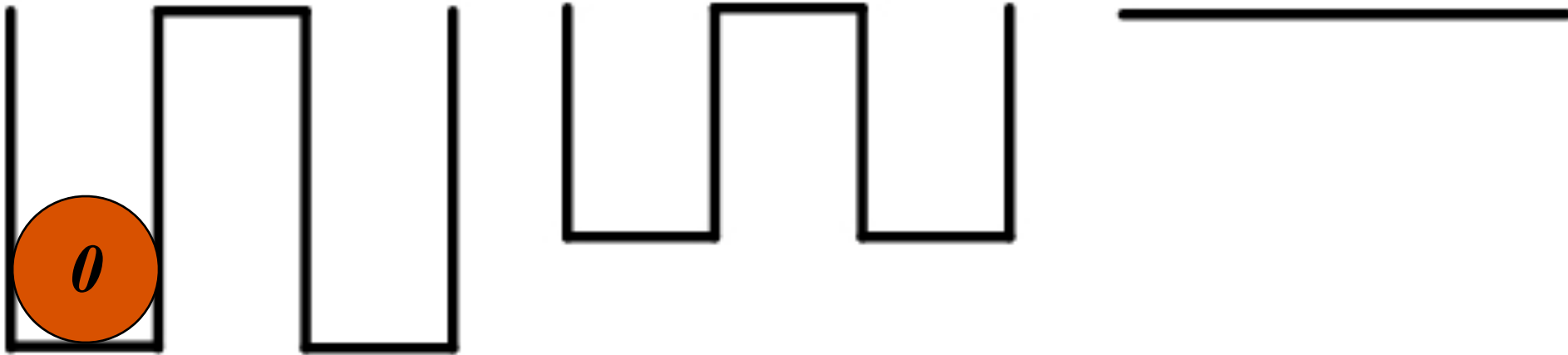
# Moving between wells

$\beta_1$

$\beta_2$

...

$\beta_c$



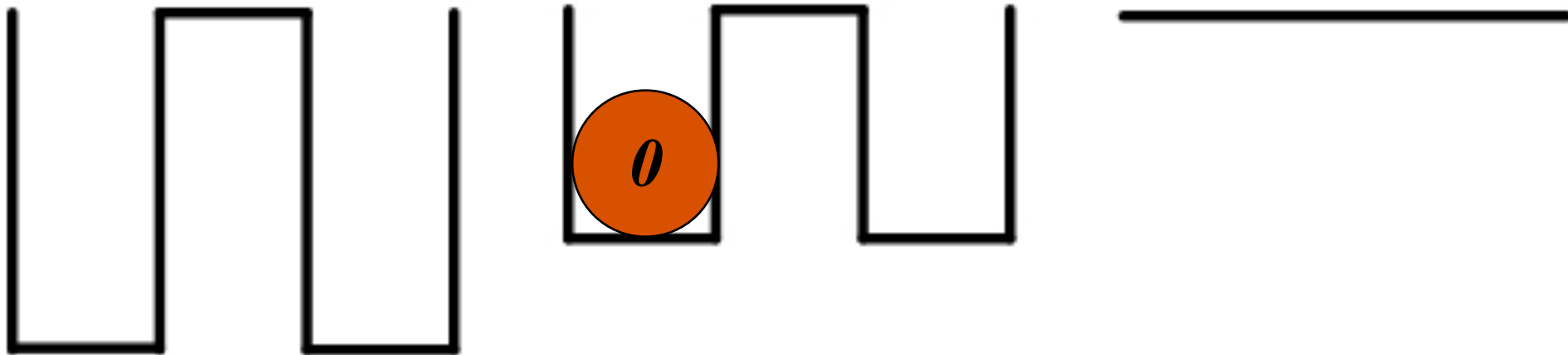
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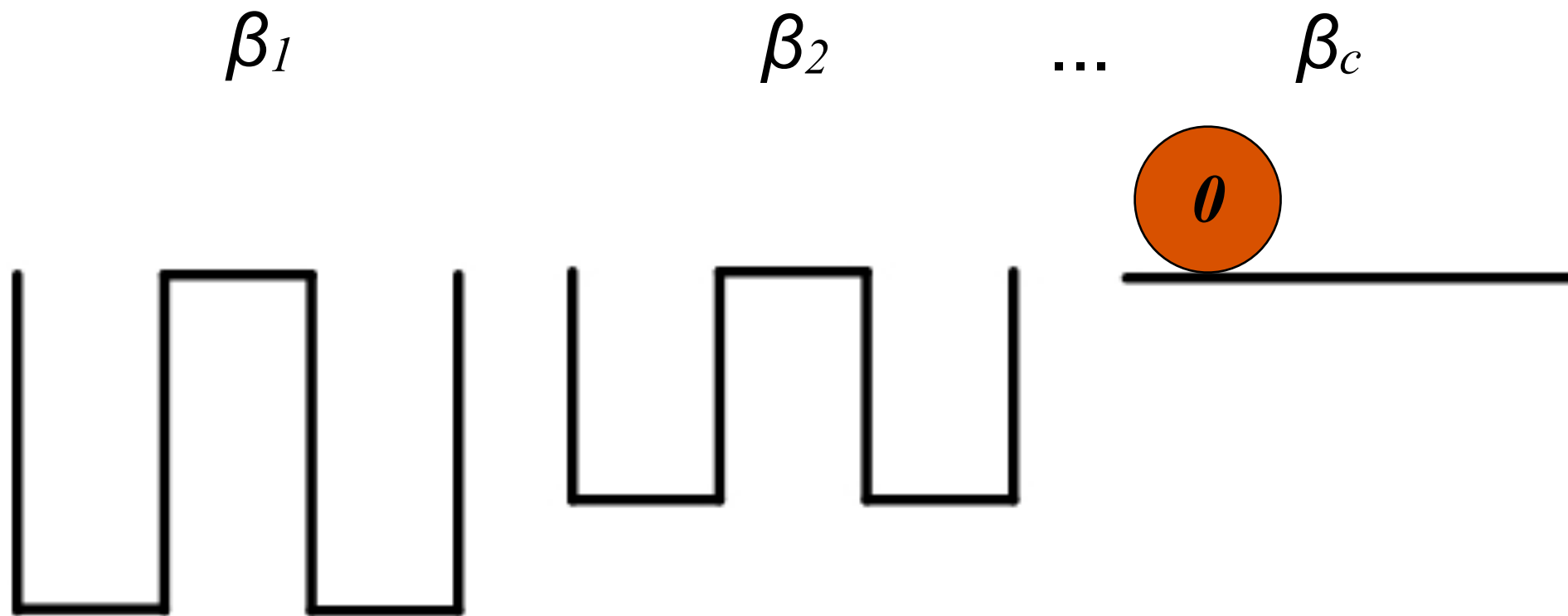
...

$\beta_c$

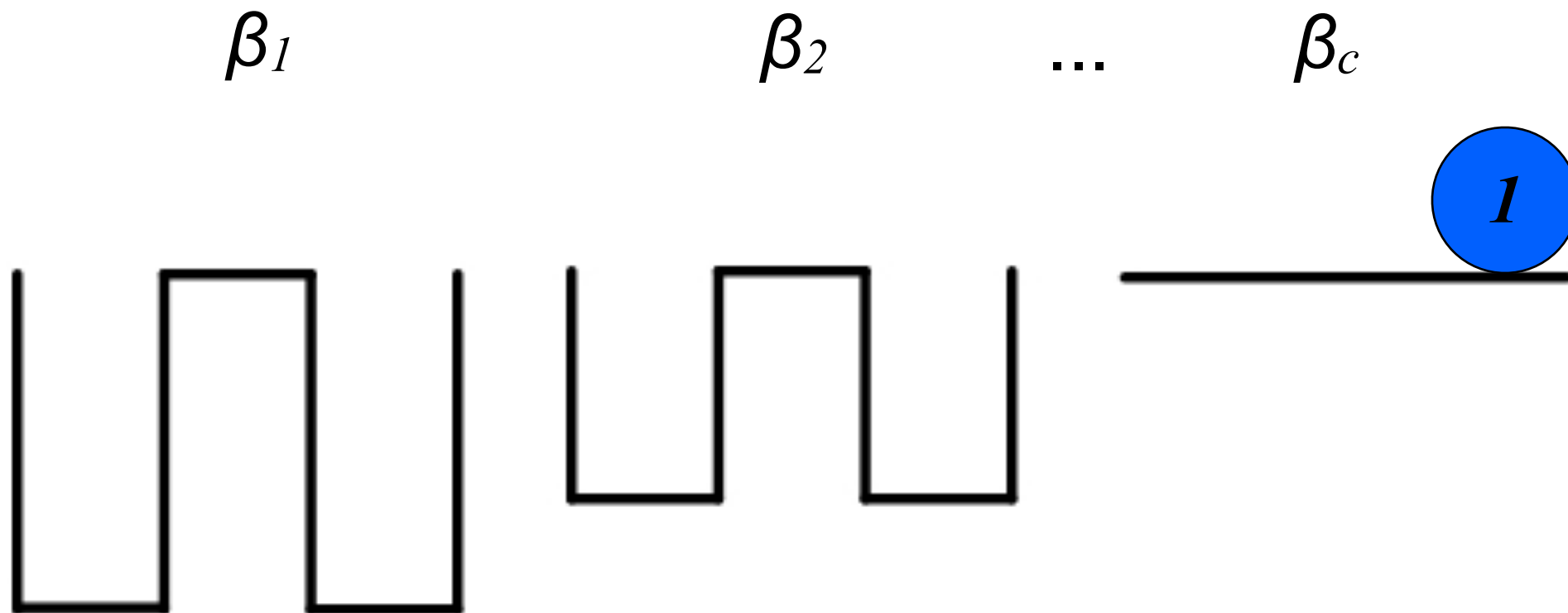




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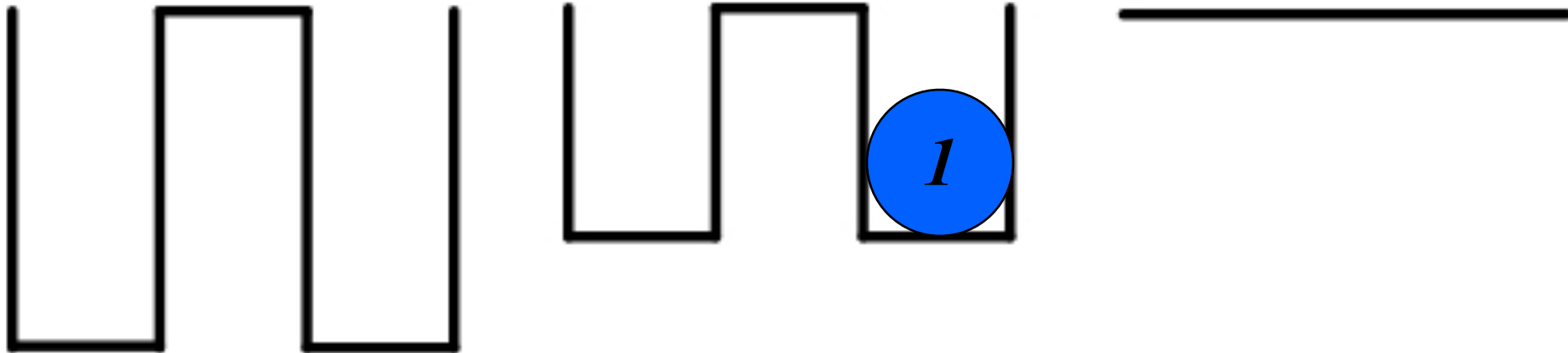
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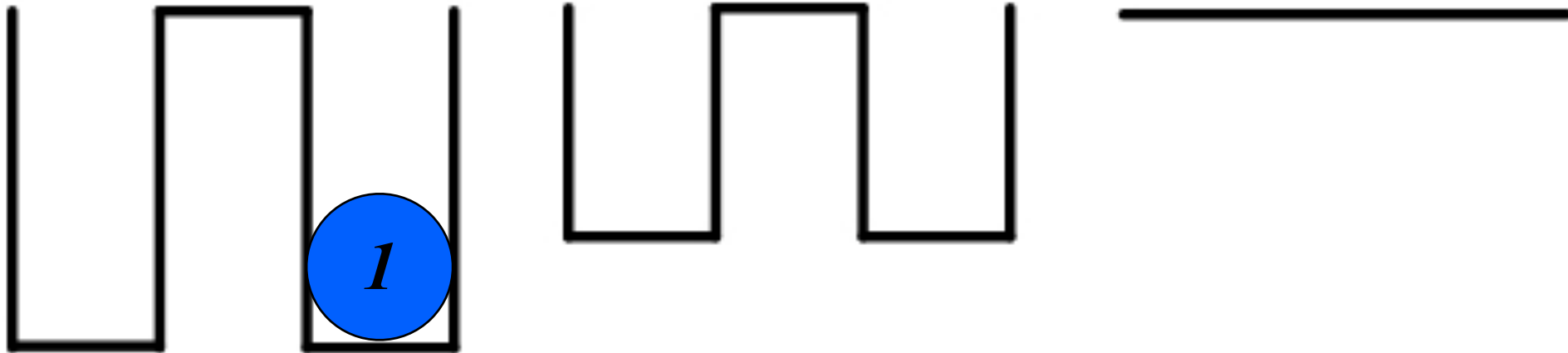
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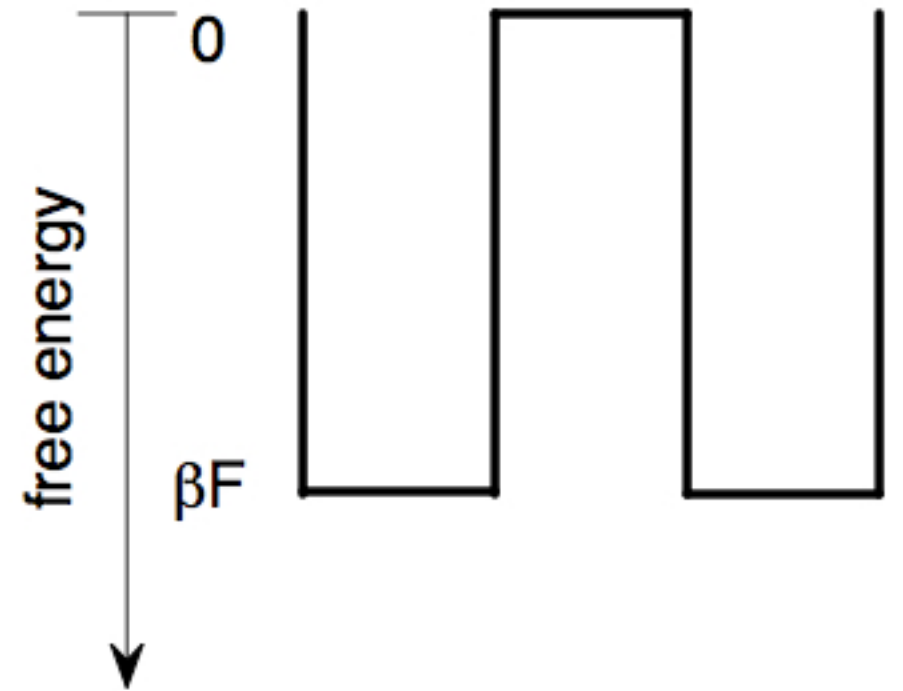
...

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# Symmetric wells ( $H=0$ )

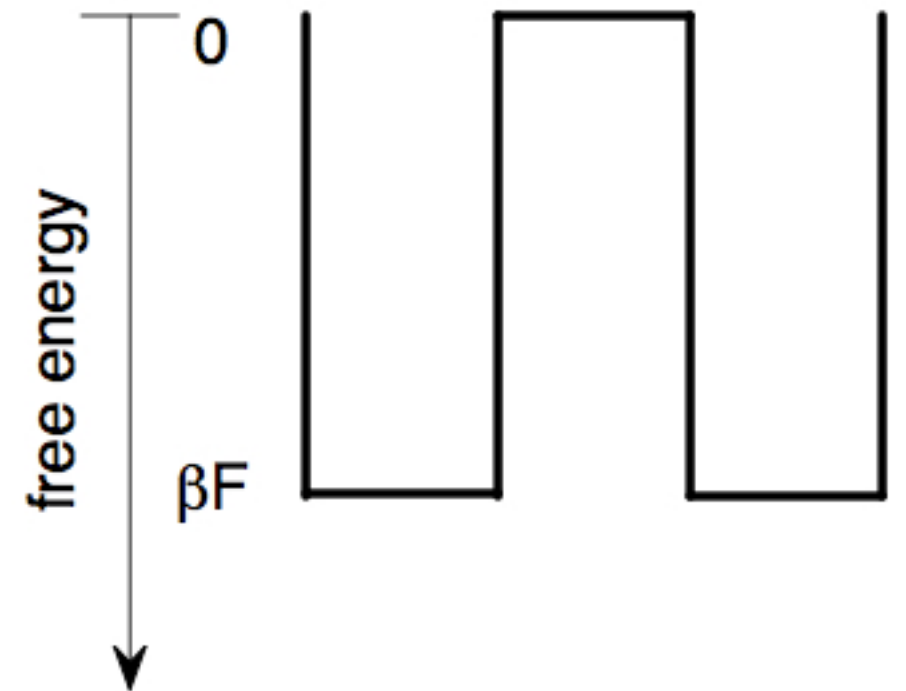
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# Symmetric wells ( $H=0$ )

- *Diffusion of replicas*

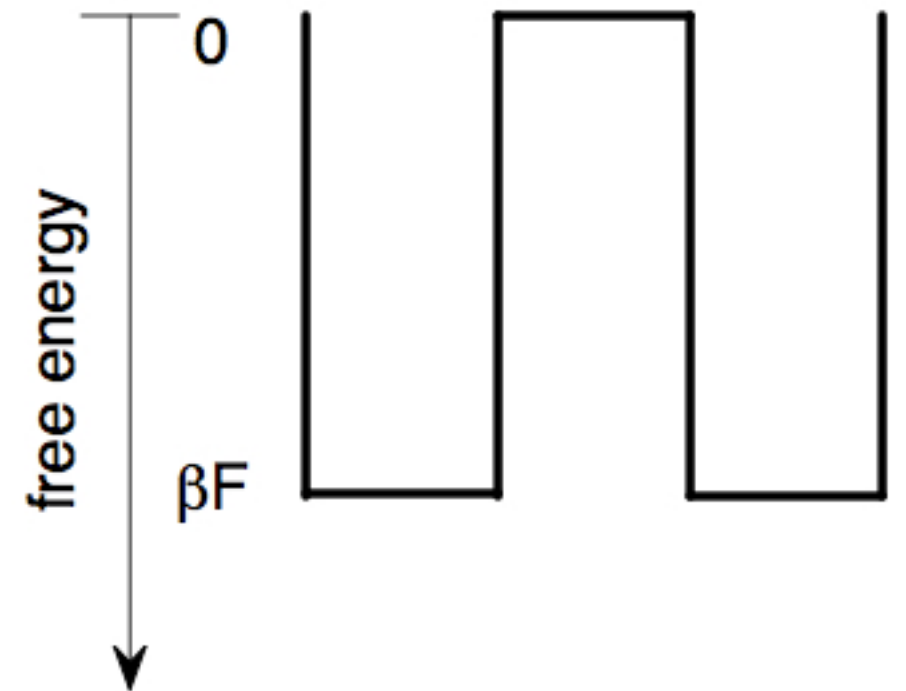
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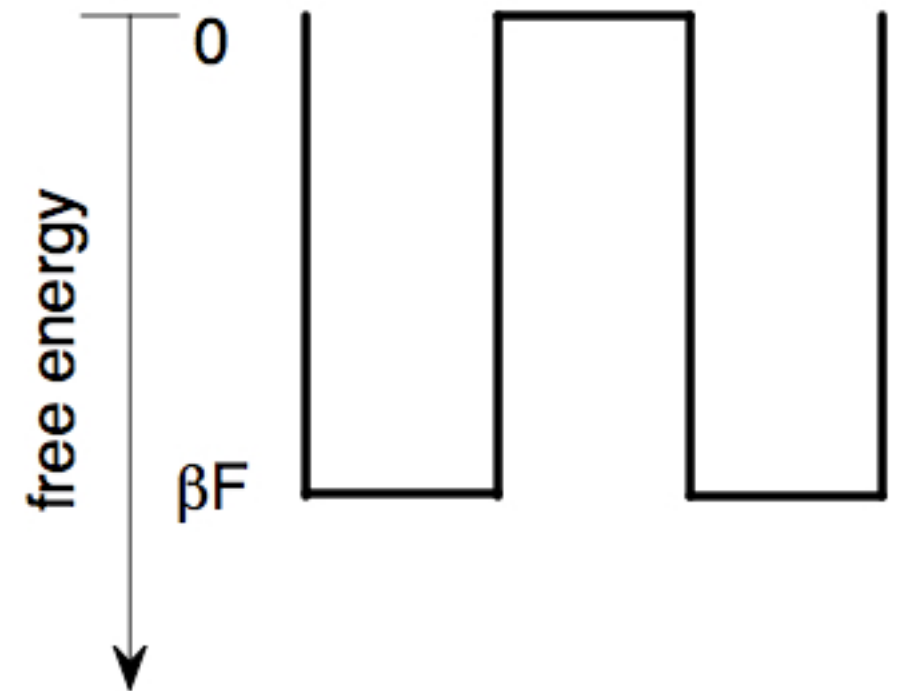
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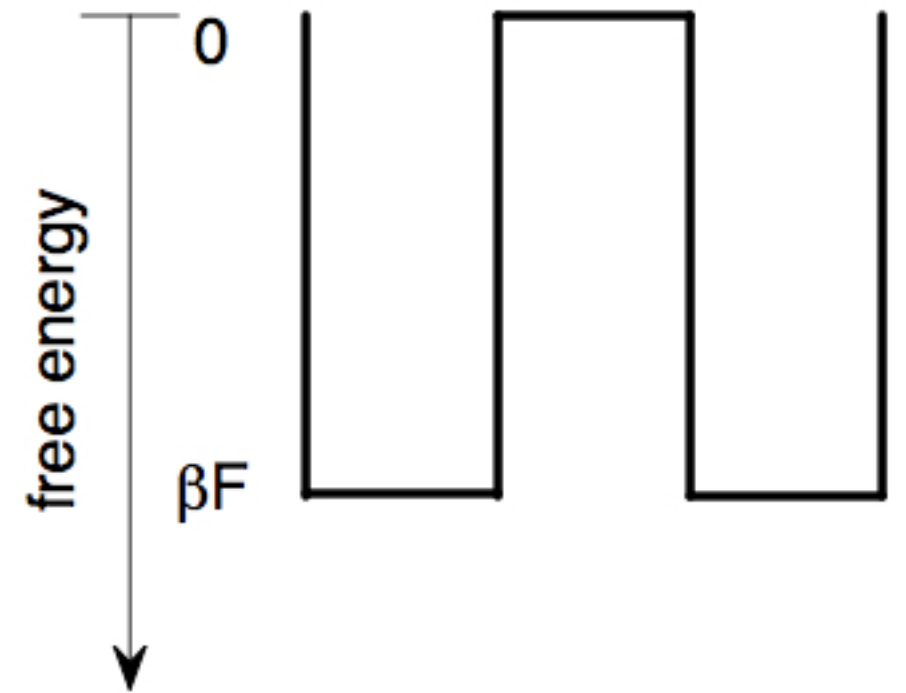




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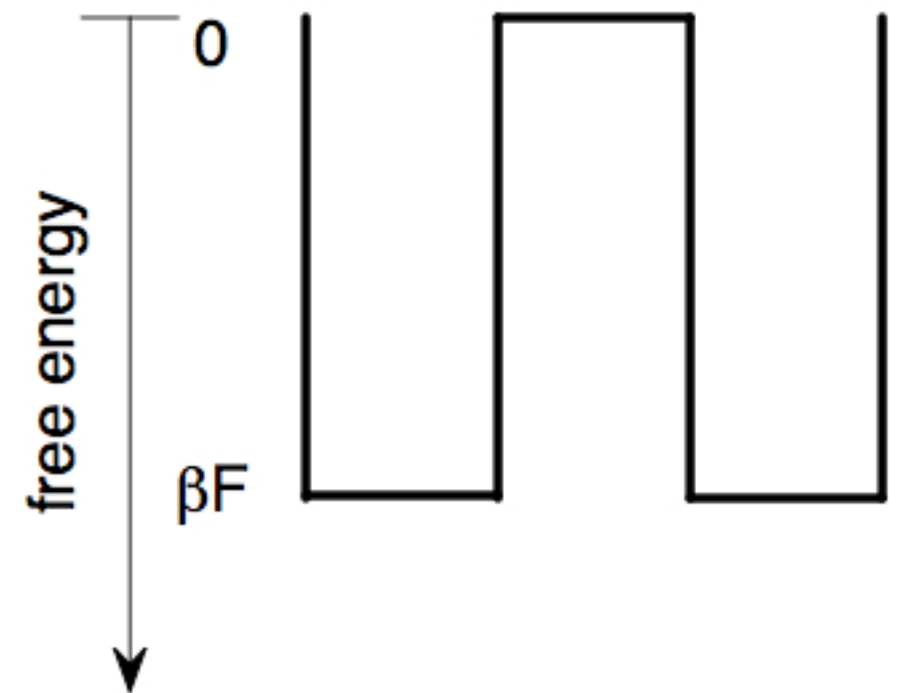
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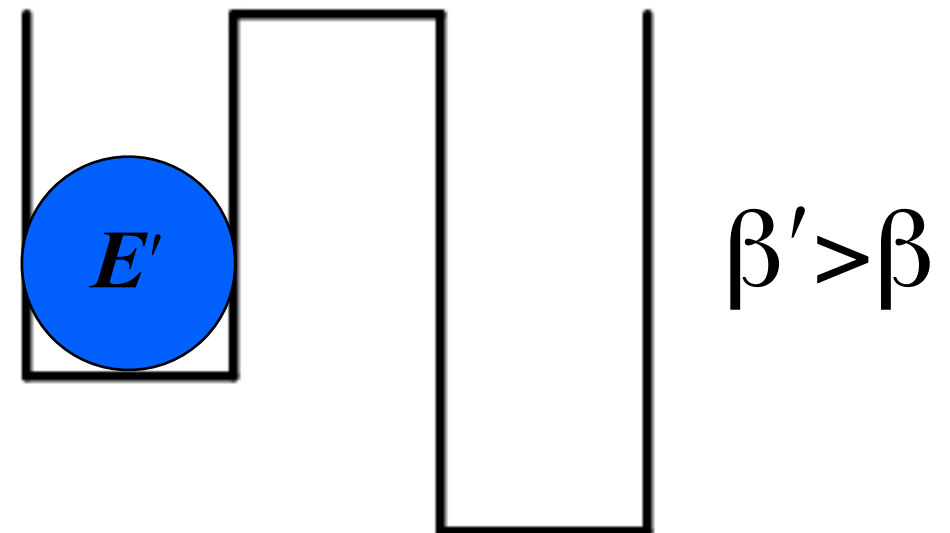
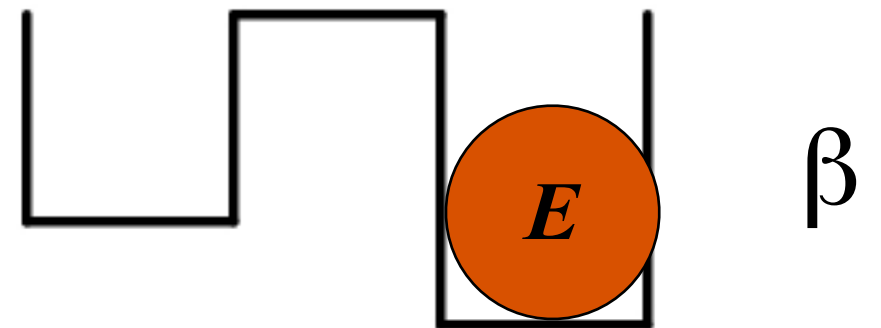
$$\tau \sim R^2 \sim \beta F$$



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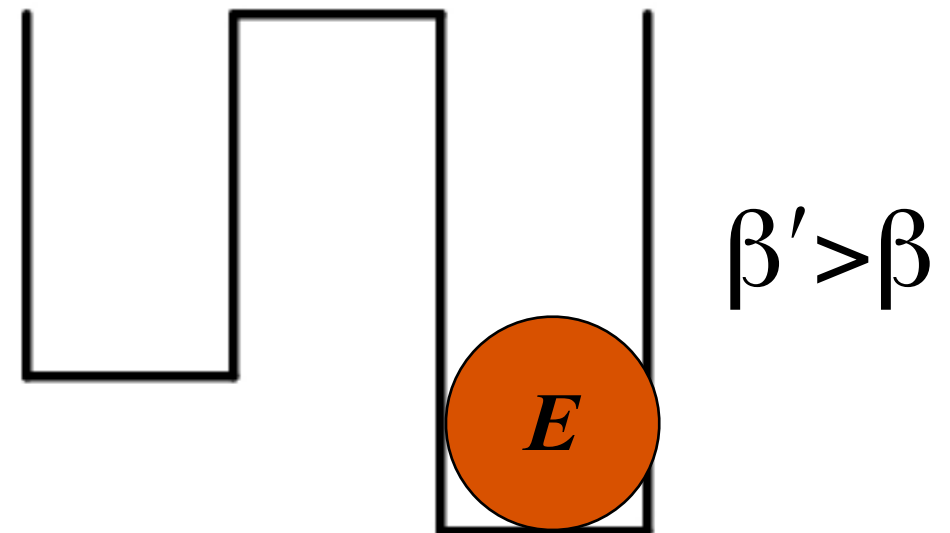
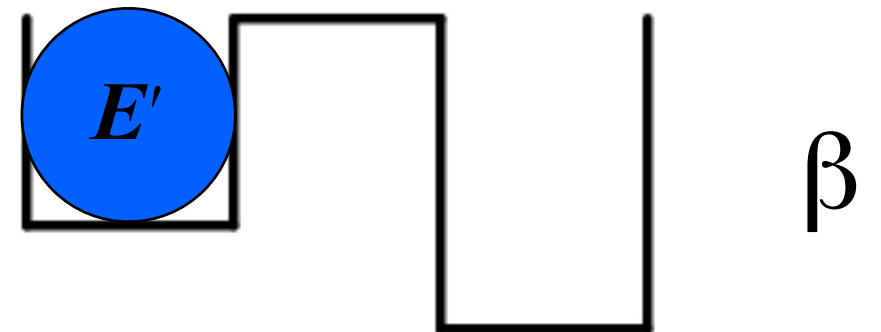
# Asymmetric wells ( $H > 0$ )

- Replica exchange is biased toward moving replicas in the deep well to lower temperature.
- *Drift and diffusion* instead of diffusion:
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- Faster equilibration than for the symmetric case!



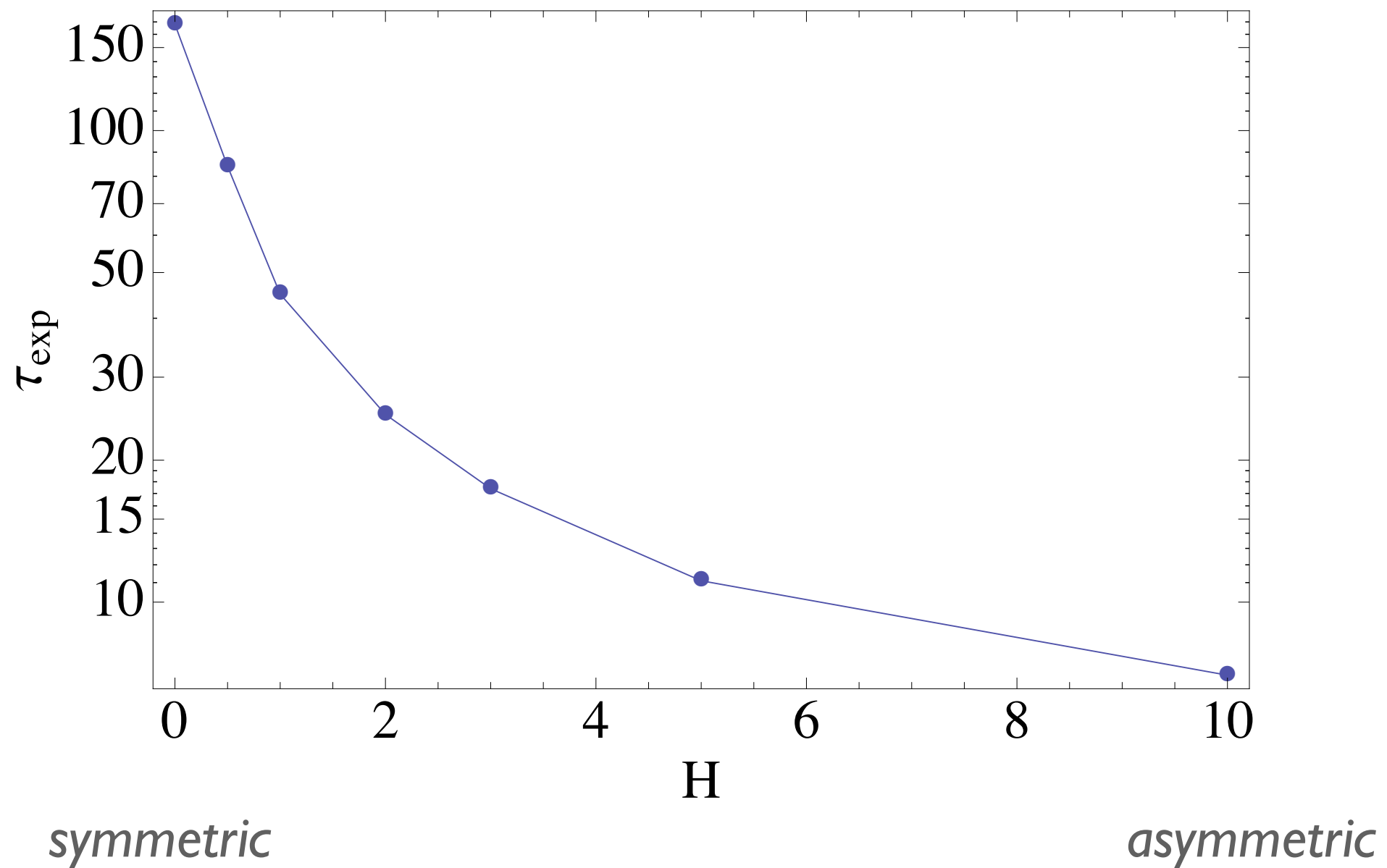
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# Equilibration time vs well asymmetry

$$\beta\delta F = -\frac{1}{2}(\beta - \beta_c)^2 H$$



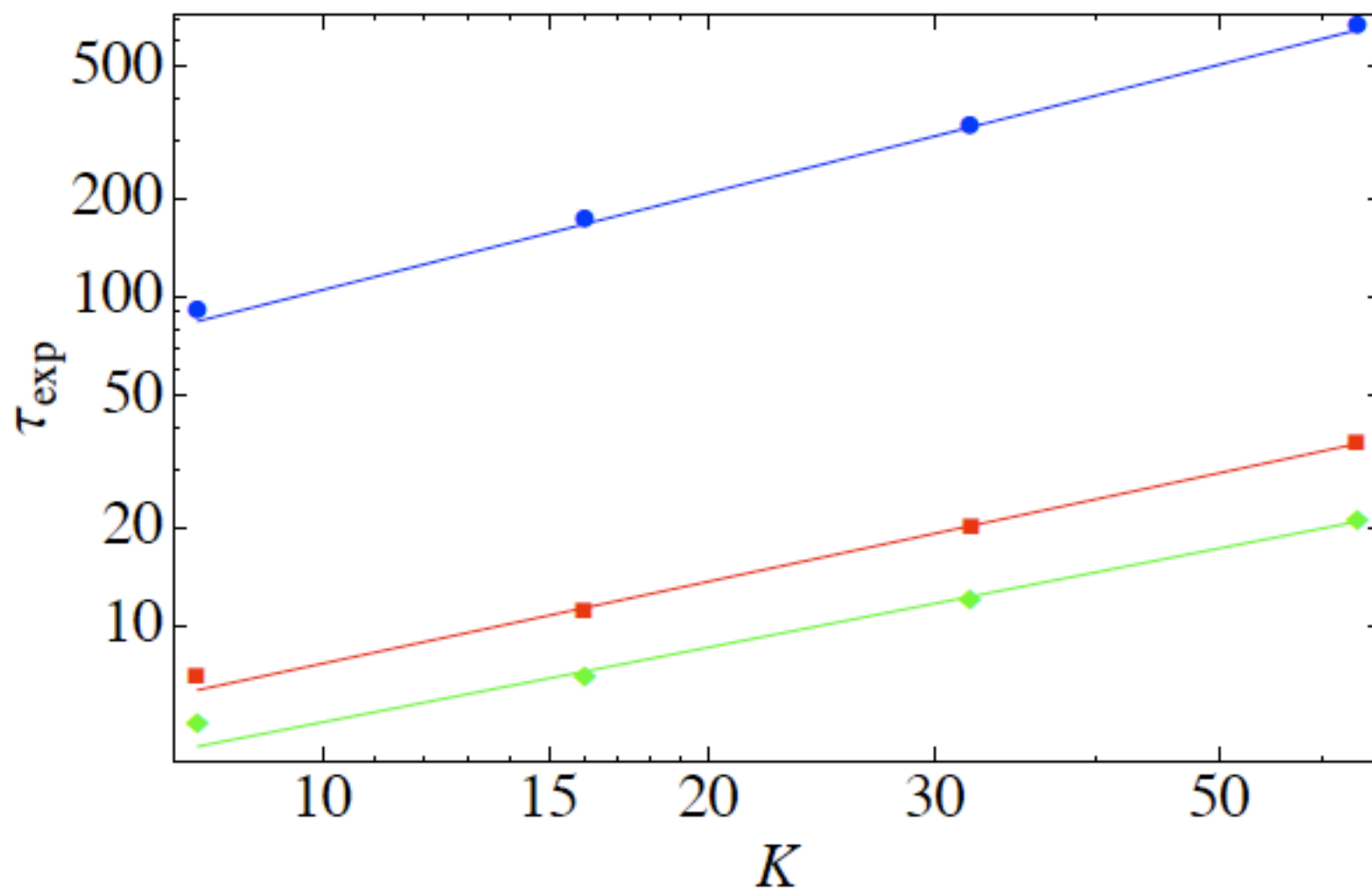
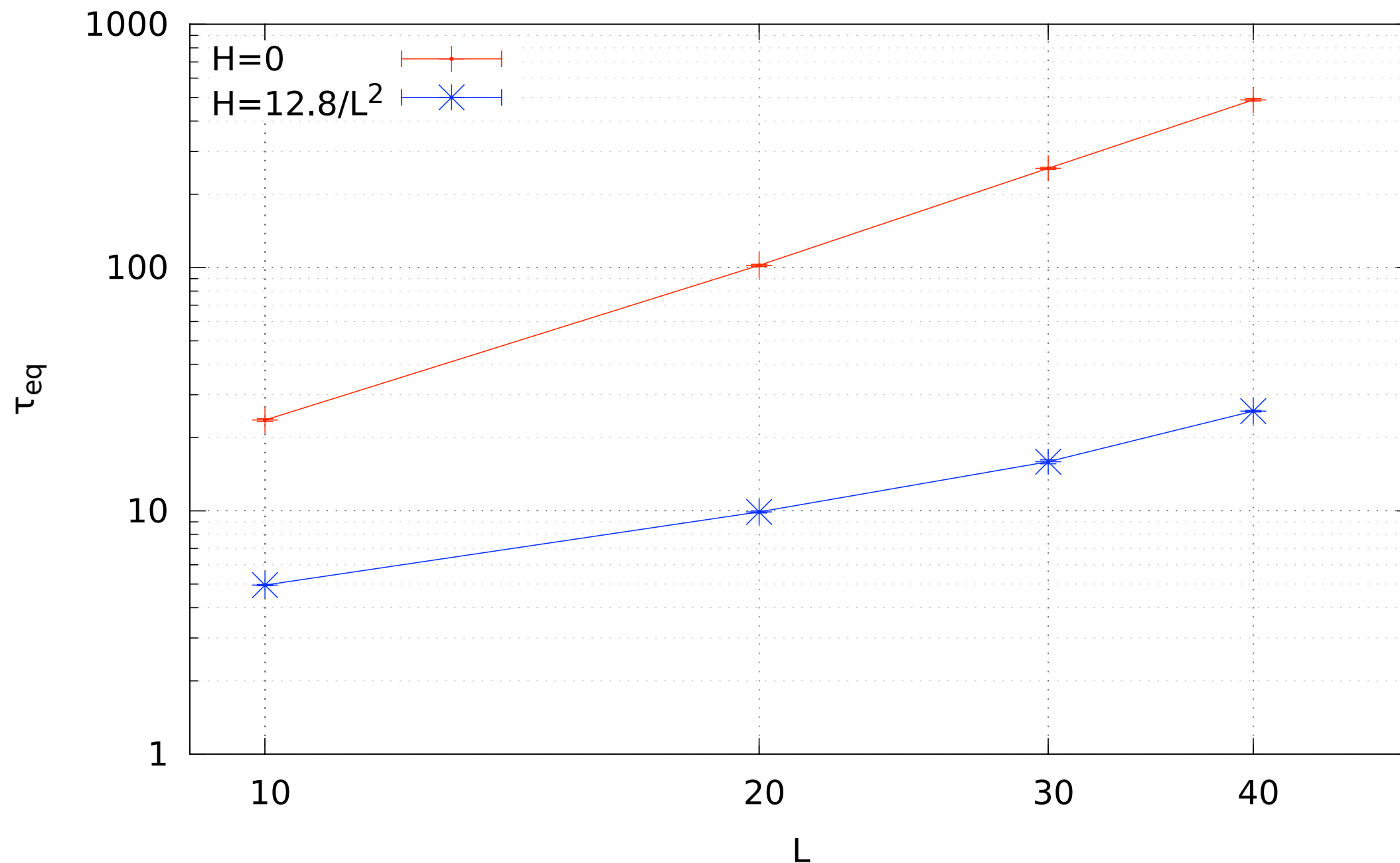


FIG. 2: The exponential autocorrelation time  $\tau_{\text{exp}}$  for the fraction of sites in the deep well vs. the well depth parameter  $K$  for  $H = 10$  (green diamonds),  $H = 5$  (red squares) and  $H = 0$  (blue circles). The lines are best power law fits,  $\tau_{\text{exp}} \sim K^x$  with  $x = 0.76, 0.83$  and  $0.99$  for  $H = 10, 5$  and  $0$ , respectively.

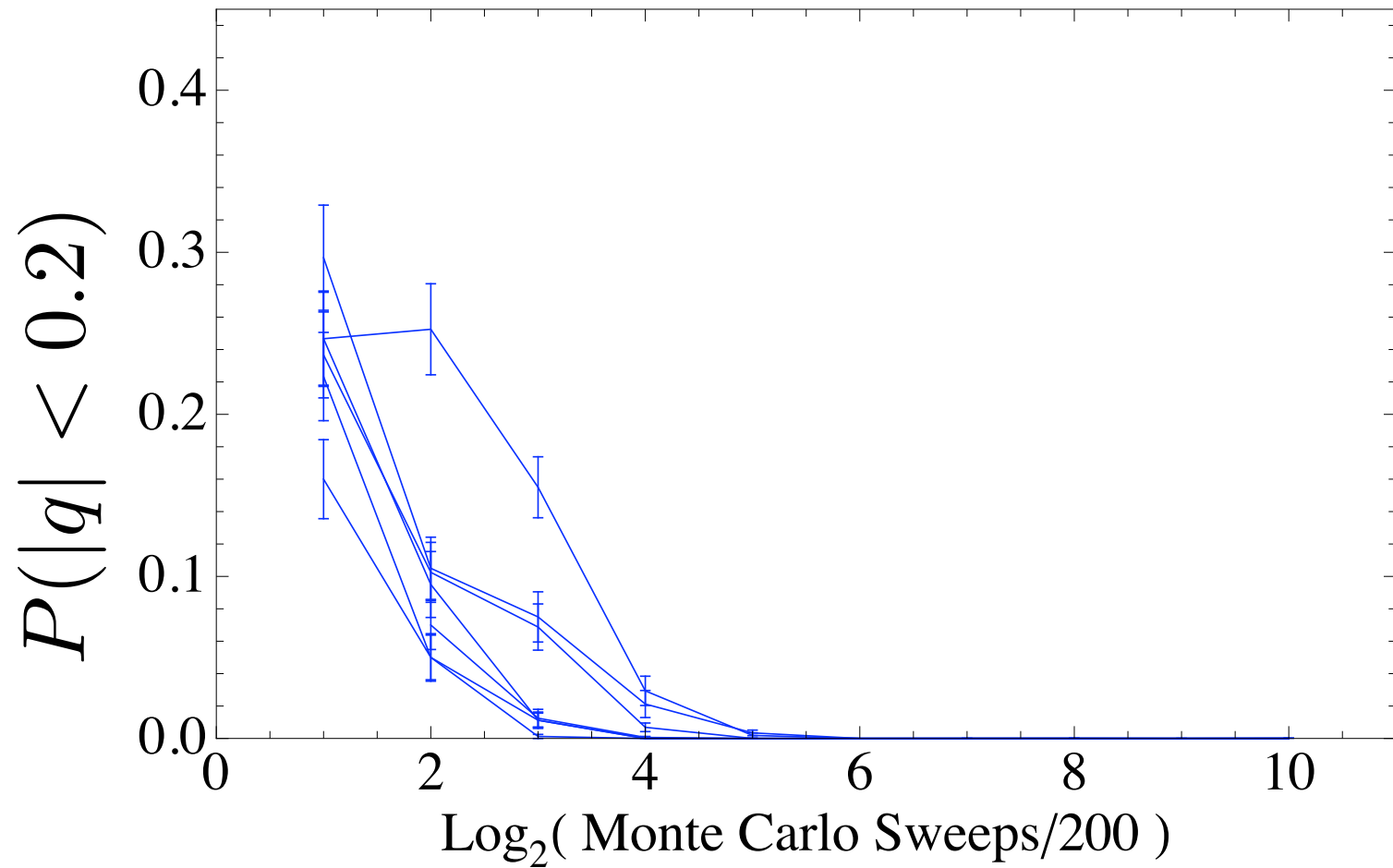
# PT for the ferromagnetic Ising model



Exponential autocorrelation time for the magnetization for the low temperature Ising model **with** and **without** a temperature dependent external field.

Ising spin glass, Gaussian disorder,  $L=8$ ,  $T/T_c=0.2$

$$P(|q| < 0.2) \approx 0$$



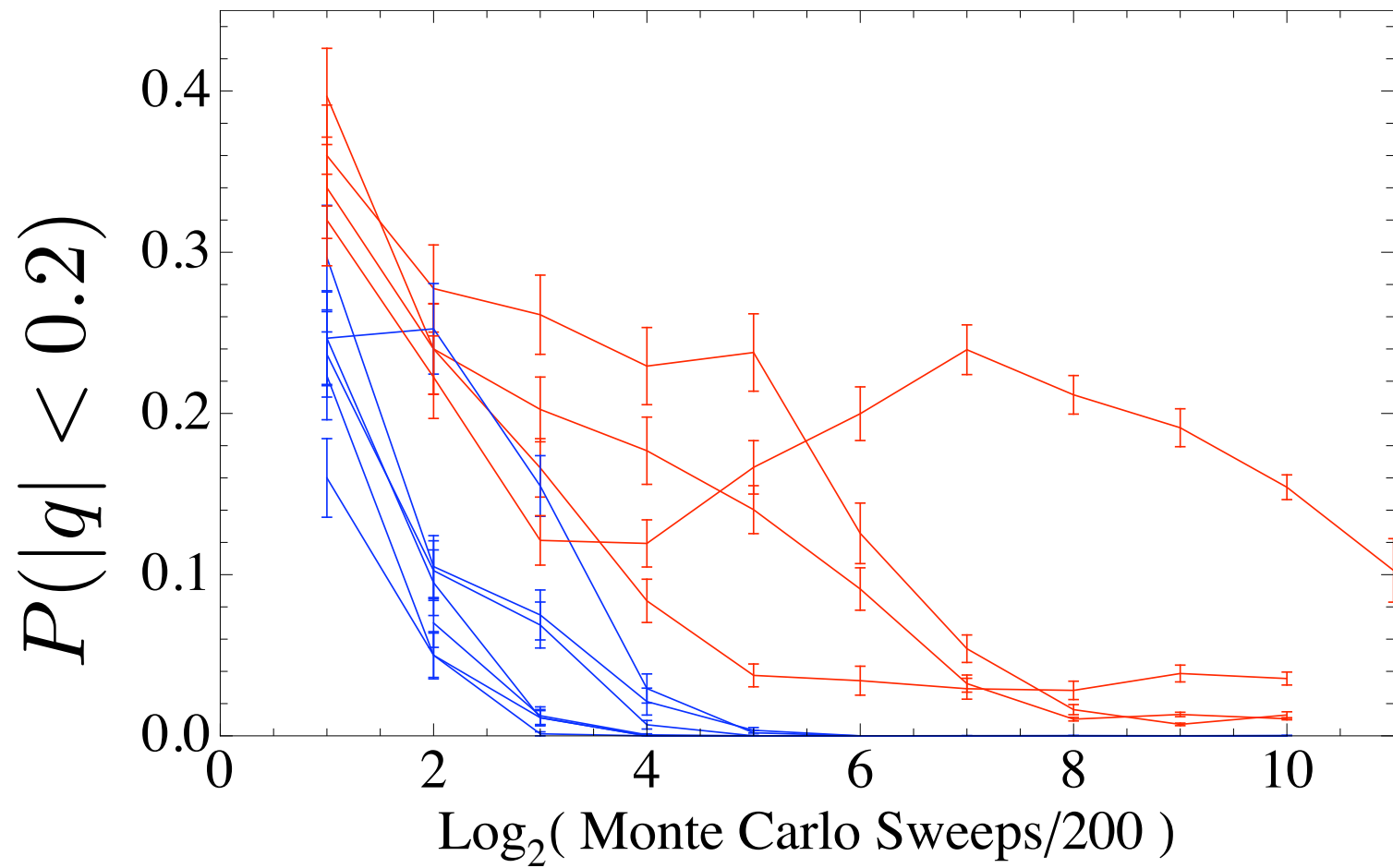


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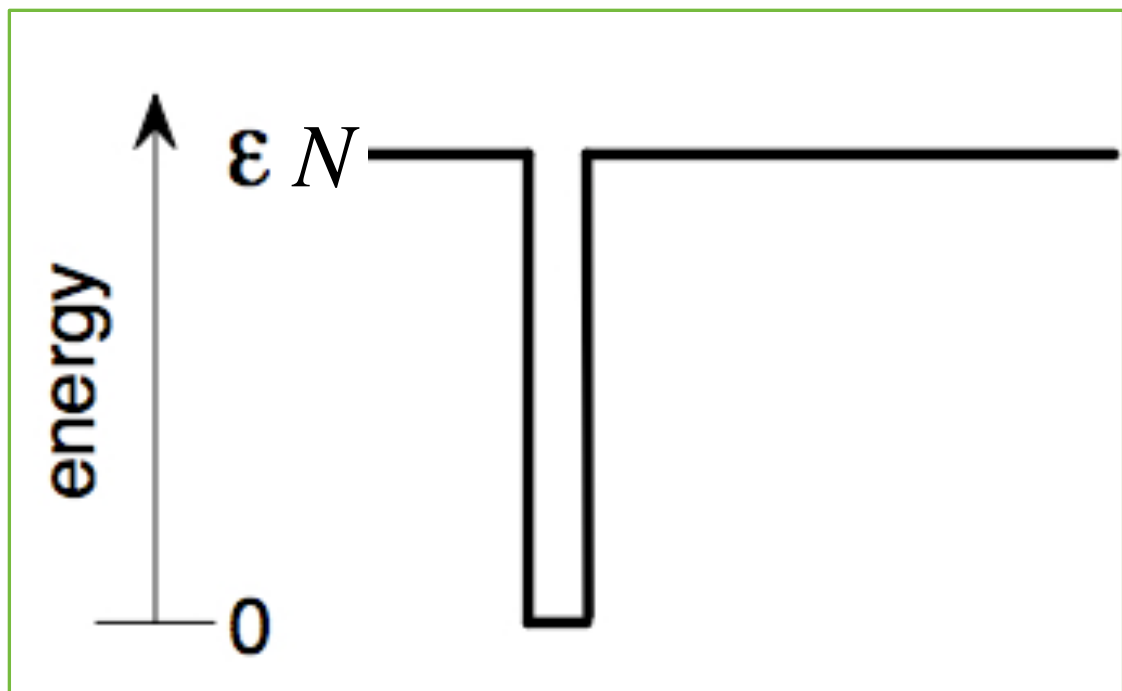


$P(|q| < 0.2) > 0$



# Weakness of Parallel Tempering

- “First-order” free energy landscapes. The simplest and worst case is the “golf-course” potential:
  - $e^N$  microstates,
  - Nearly all have energy  $\varepsilon N > 0$
  - A small fraction,  $e^{-\varepsilon\beta_c N}$  are ground states with energy zero. There is a pseudo phase transition at  $\beta_c$
  - Probability to be in a ground state:



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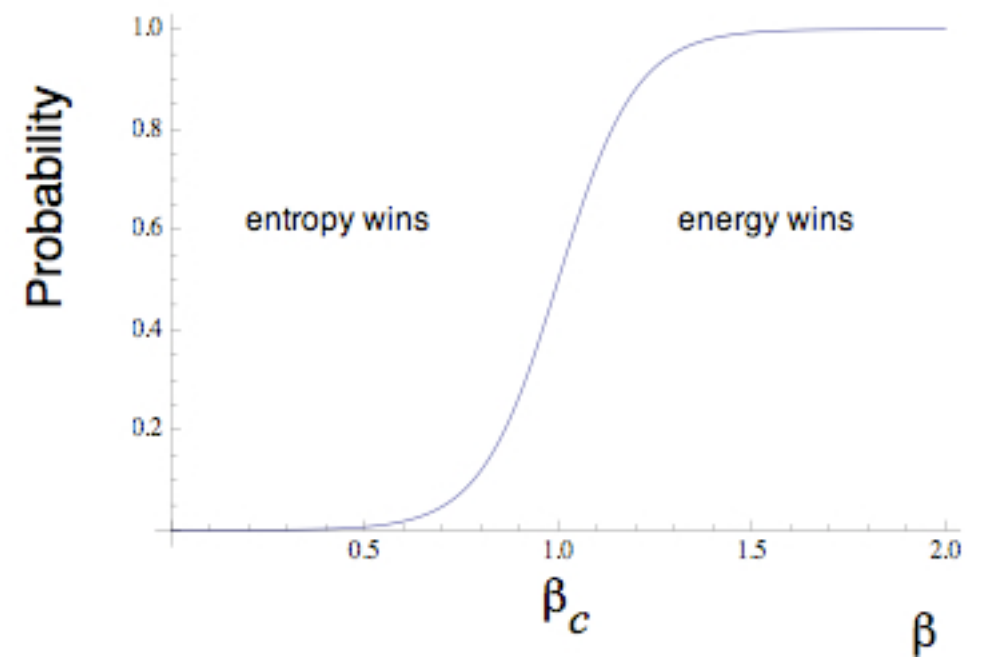
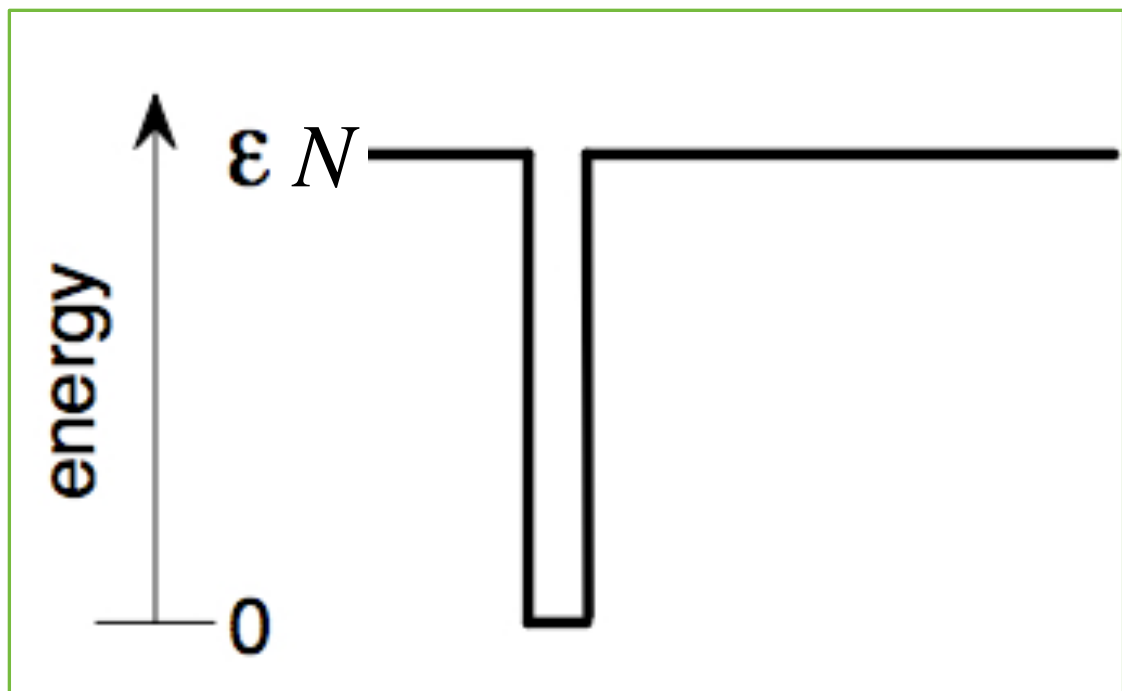
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- Probability to be in a ground state:

$$\frac{1}{1 + e^{-(\beta - \beta_c)\epsilon N}}$$



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- *Before equilibration round trip-time is misleading.*

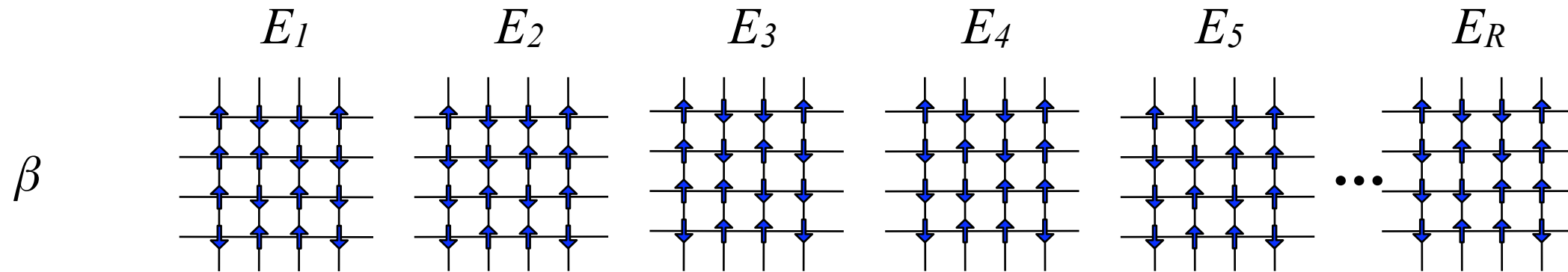


# Population Annealing

K. Hukushima and Y. Iba, in *THE MONTE CARLO METHOD IN THE PHYSICAL SCIENCES: Celebrating the 50th Anniversary of the Metropolis Algorithm*, edited by J. E. Gubernatis (AIP, 2003), vol. 690, pp. 200–206.

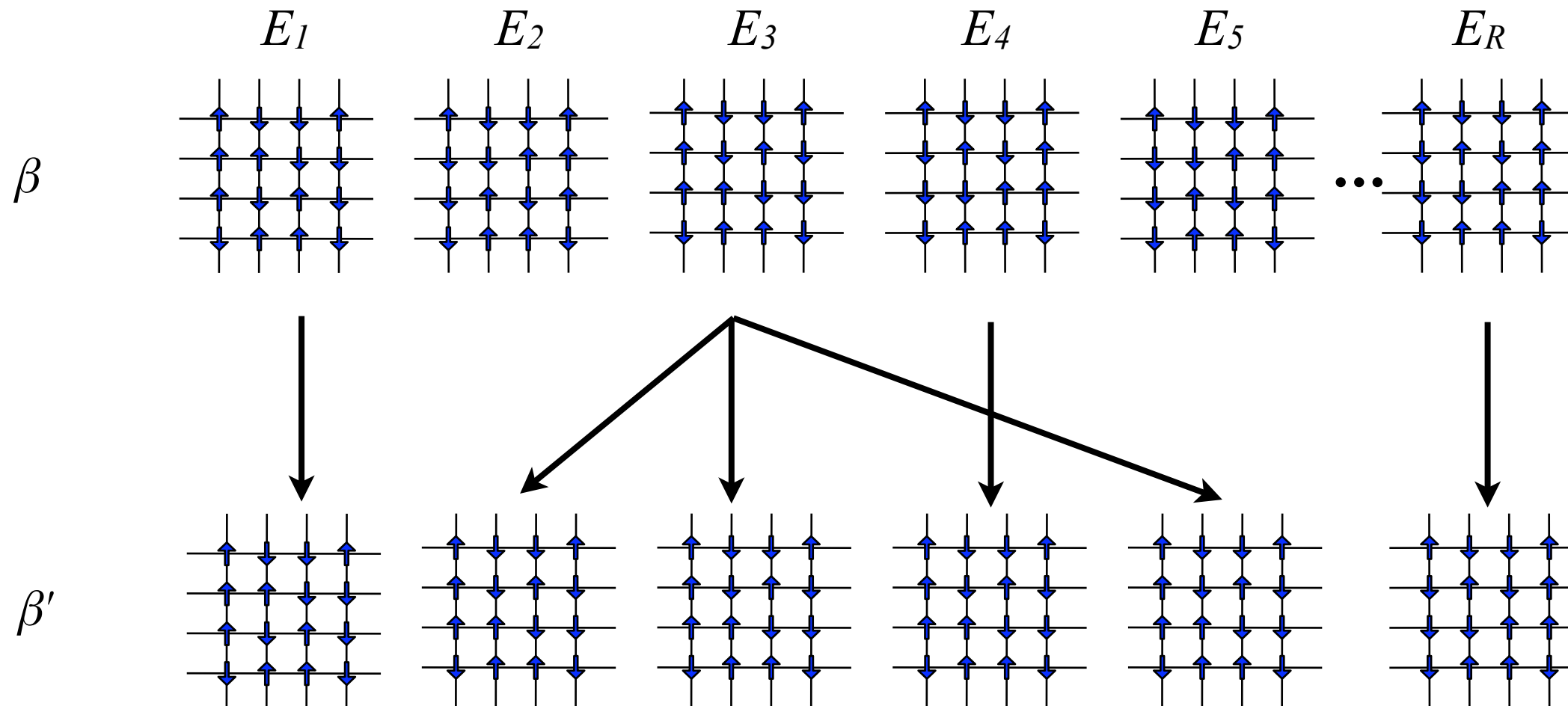
- Modification of simulated annealing for equilibrium sampling.
- A population of replicas is cooled according to an annealing schedule. Each replica is acted on by a MCMC (e.g. Metropolis) at the current temperature.
- During each temperature step, the population is differentially resampled according to Boltzmann weights to maintain equilibrium.

# Population Annealing



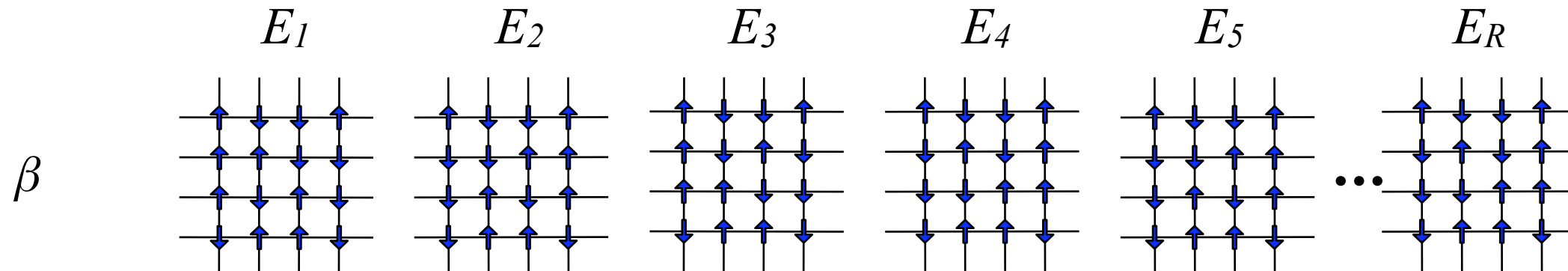
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# Population Annealing



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# Temperature Step

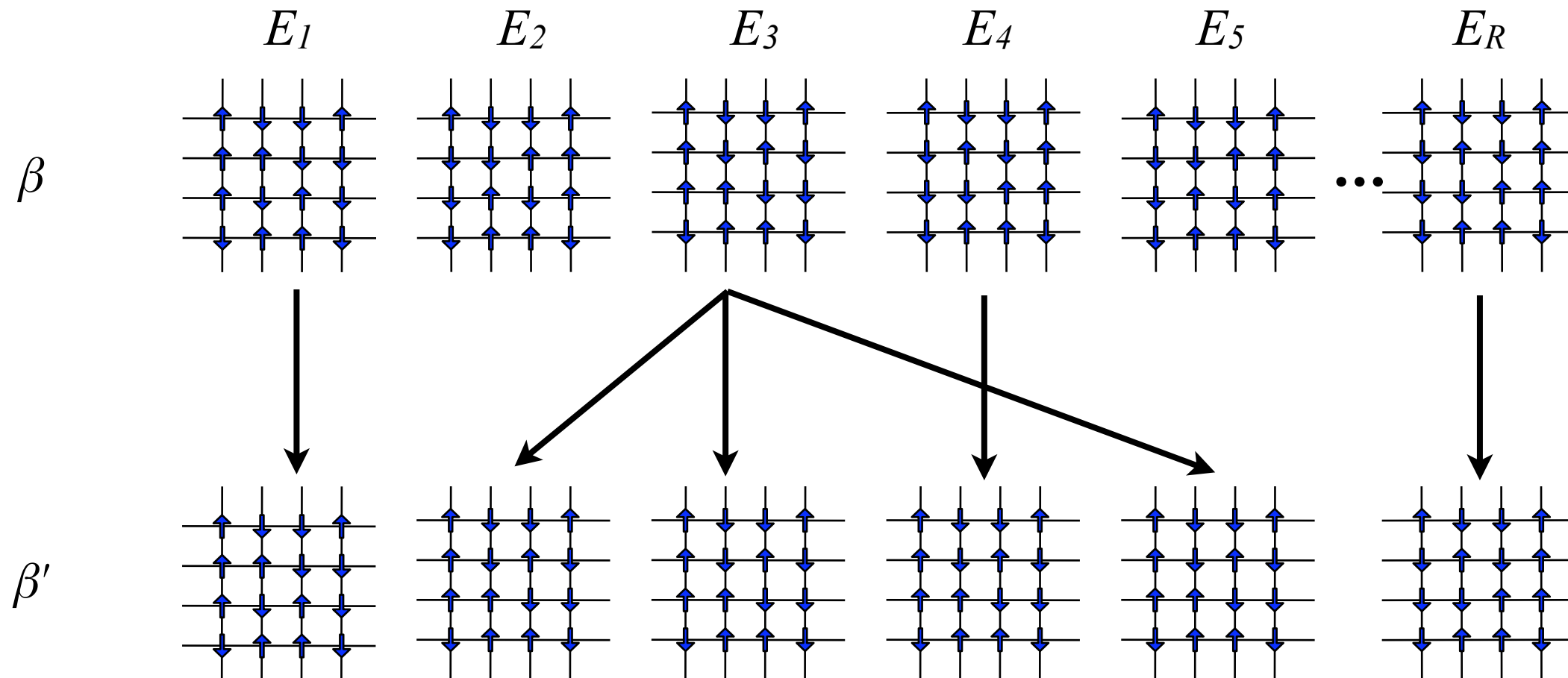


$$\tau_j(\beta, \beta') = \frac{\exp [-(\beta' - \beta)E_j]}{Q(\beta, \beta')}$$

$$Q(\beta, \beta') = \frac{\sum_{j=1}^R \exp [-(\beta' - \beta)E_j]}{R}$$

Replica  $j$  is reproduced  $n_j$  times where  $n_j$  is a Poisson random number with mean  $\tau_j$ .

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# Strengths of Population Annealing

- Massively parallel
- Direct measurement of the free energy

$$-\beta_k F(\beta_k) = \sum_{\ell=k}^K \ln Q(\beta_{\ell+1}, \beta_{\ell}) + \beta_K F(\beta_K) \quad Q(\beta, \beta') = \frac{\sum_{j=1}^R \exp[-(\beta' - \beta)E_j]}{R}$$

- Comparable efficiency to parallel tempering for Ising spin glasses.

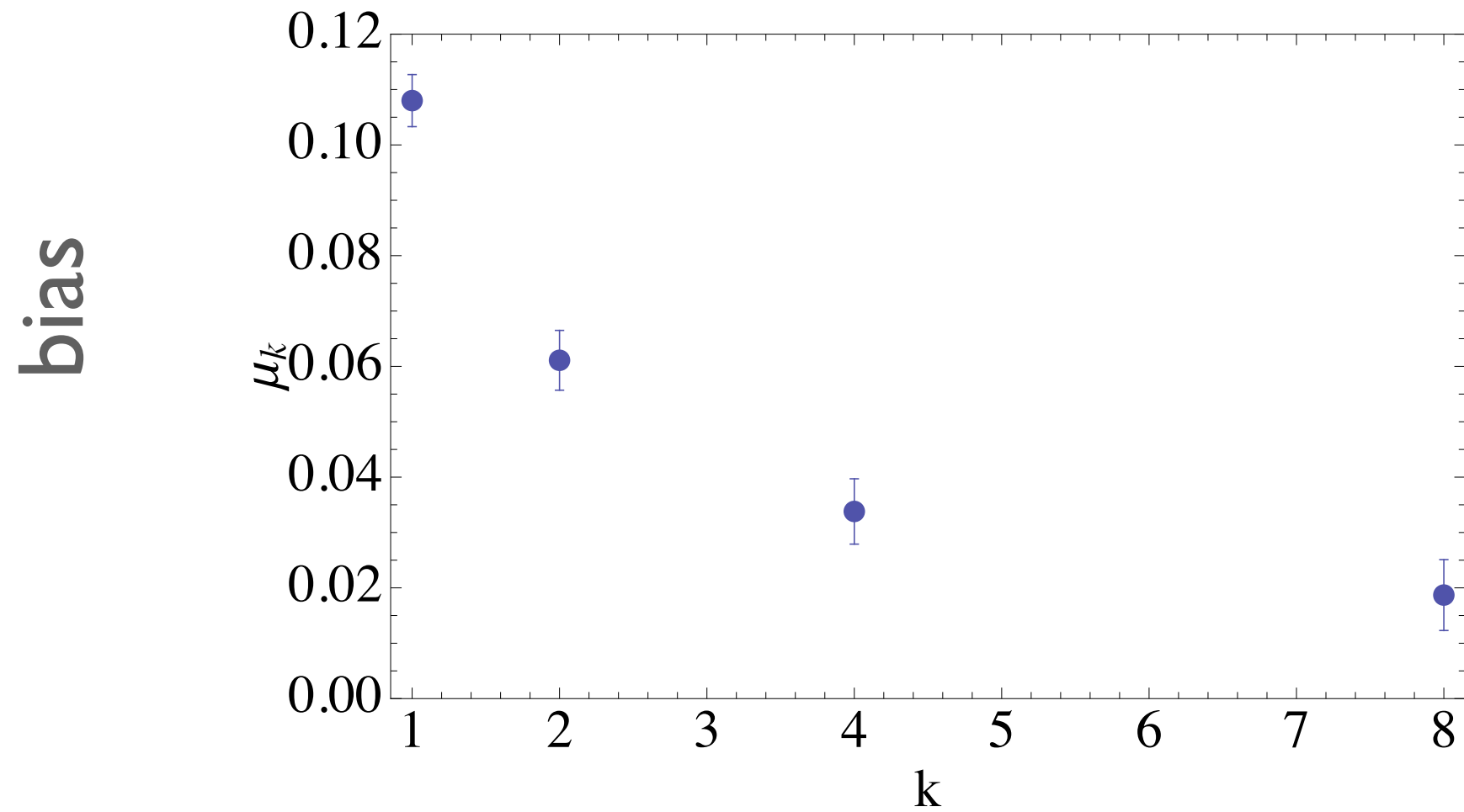
# Weighted Averages

JM, arXive:1006.0252

- Results using small number of replicas are biased.
- Results from independent runs can be combined and biases reduced using weighted averages.
- Observables from each run weighted by the exponential of the free energy estimator for that run.

$$\bar{A}(\beta) = \sum_{m=1}^M \tilde{A}_m(\beta) \omega_m(\beta) \quad \omega_m(\beta) = \frac{e^{-\beta \tilde{F}_m(\beta)}}{\sum_{i=1}^M e^{-\beta \tilde{F}_i(\beta)}}$$

- Number of runs needed for unbiased results determined by variance of the free energy estimator--an intrinsic measure of equilibration.



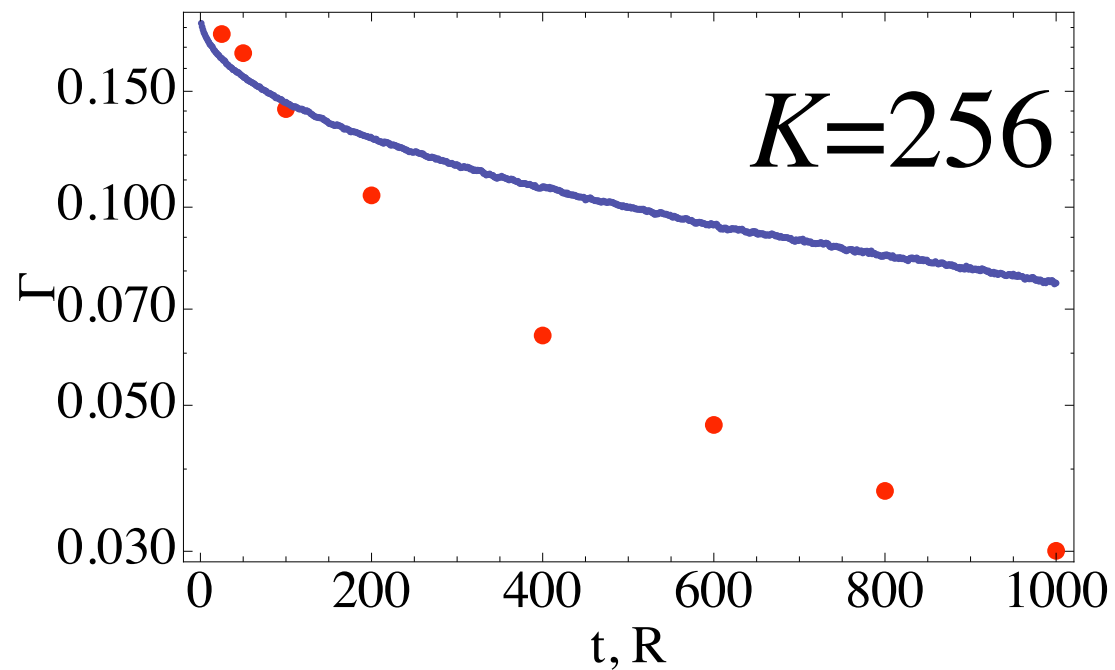
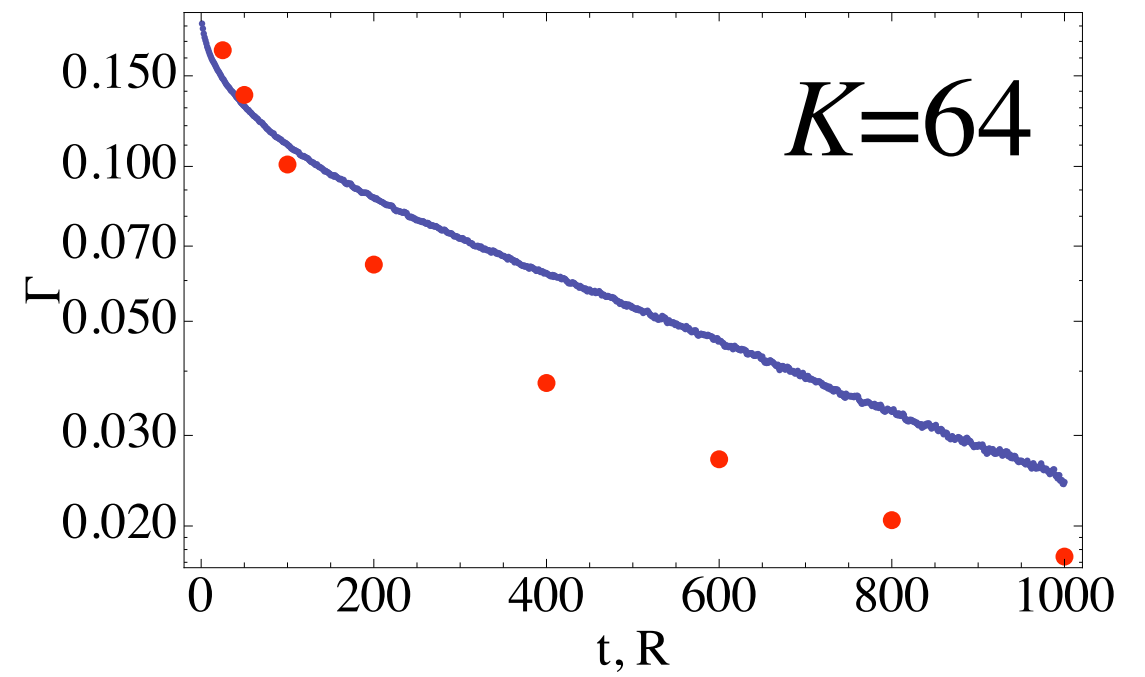
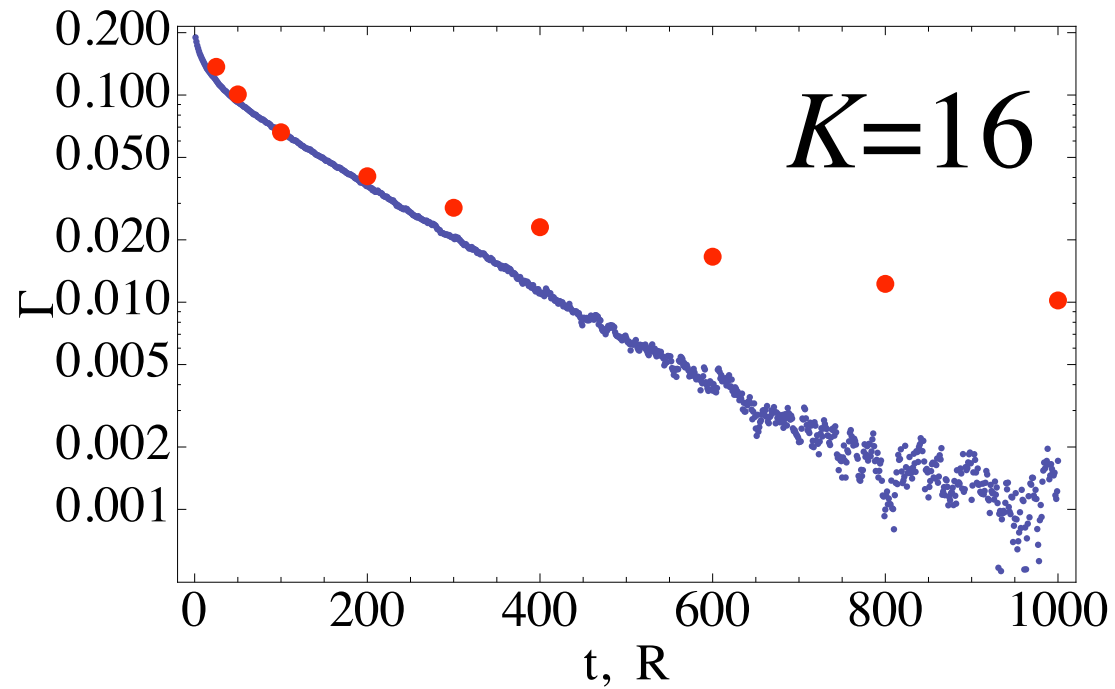
number of runs

Bias in the energy for the 1D Ising spin glass as a function of ensemble size.



<b>Population Annealing</b>	<b>Parallel Tempering</b>
Sequential Monte Carlo	Markov Chain Monte Carlo
# Replicas ( $R$ ) <i>space</i>	$\approx$ #Sweeps ( $t$ ) <i>parallel time</i>
#Temperature steps ( $T$ ) <i>parallel time</i>	$\approx$ #Replicas ( $R$ ) <i>space</i>

# Compare PA and PT for the Two-Well Model



● = PA  
— = PT

$H=0.1$

$$\Gamma = \langle \text{Prob} [\sigma = +1] \rangle - \langle \text{Prob} [\sigma = +1] \rangle_{\text{eq}}$$

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#Temperature steps ( $T$ ) <i>parallel time</i>	#Replicas ( $R$ ) <i>space</i>
$(A - A_{eq}) \approx R_0/R$	$(A - A_{eq}) \approx e^{(-t/\tau)}$
$R_0 \sim K^{1/2}$ $T \sim K^{1/2}$	$\tau \sim K$ $R \sim K^{1/2}$

# Conclusions

- Both parallel tempering and population annealing are efficient at finding equilibria between free energy minima with large basins of attraction separated by large barriers.
  - For PT efficiency is greatest when the minima are highly asymmetric
- Neither method is useful, except via brute force parallelism, for the case of free energy minima with small basins of attraction.
- Population annealing is comparably efficient to parallel tempering and has several features to recommend it:
  - Highly parallel
  - Direct measurement of free energy
  - Built-in equilibration criteria (variance of the free energy estimator)